

# HEAT KERNELS AND GREEN'S FUNCTIONS ON LIMIT SPACES

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ABSTRACT. In this paper, we study the behavior of the Laplacian on a sequence of manifolds  $\{M_i^n\}$  with a lower bound in Ricci curvature that converges to a metric-measure space  $M_\infty$ . We prove that the heat kernels and Green's functions on  $M_i^n$  will converge to some integral kernels on  $M_\infty$  which can be interpreted, in different cases, as the heat kernel and Green's function on  $M_\infty$ . We also study the Laplacian on noncollapsed metric cones; these provide a unified treatment of the asymptotic behavior of heat kernels and Green's functions on noncompact manifolds with nonnegative Ricci curvature and Euclidean volume growth. In particular, we get a unified proof of the asymptotic formulae of Colding-Minicozzi, Li and Li-Tam-Wang.

## INTRODUCTION

Assume  $M^n$  is an  $n$  dimensional Riemannian manifold with a lower bound in Ricci curvature,

$$(0.1) \quad \text{Ric}_{M^n} \geq -(n-1)\Lambda,$$

where  $\Lambda \geq 0$ . By the Bishop-Gromov inequality, we have a *uniform* volume doubling condition,

$$(0.2) \quad \text{Vol}(B_{2R}(p)) \leq 2^\kappa \text{Vol}(B_R(p)),$$

here we can take  $\kappa = n$  if  $\Lambda = 0$ ; if  $\Lambda > 0$ , we require that  $R$  is bounded from above, say,  $R < D$  for some  $D > 0$ . Moreover, there is a *uniform* Poincare inequality

$$(0.3) \quad \frac{1}{\text{Vol}(B_R(p))} \int_{B_R(p)} |f - f_{p,R}| \leq \tau R \left( \frac{1}{\text{Vol}(B_R(p))} \int_{B_R(p)} |df|^2 \right)^{\frac{1}{2}},$$

where  $f_{p,R}$  is the average of  $f$  on  $B_R(p)$ ; see [Bu], [Ch3], [HaKo] and the references therein.

We assume, throughout this paper,  $\{M_i^n\}$  is a sequence of complete Riemannian manifolds with (0.1) that converges in the pointed measured Gromov-Hausdorff sense, to a metric space  $M_\infty$ ; we write  $M_i^n \xrightarrow{d_{GH}} M_\infty$ ,  $d_{GH}$  is the Gromov-Hausdorff distance. In particular, (0.2) holds on  $M_\infty$ . One can show that (0.3) and the *segment inequality*, which is stronger than (0.3), hold on  $M_\infty$  as well; note on  $M_\infty$ , the role of  $|du|$  in (0.3) is played by  $g_f$ , the *minimal generalized upper gradient*, see [Ch3], [ChCo4].

In Cheeger's paper [Ch3], a significant portion of analysis on smooth manifolds was extended to metric-measure spaces satisfying (0.2), (0.3). In [Ch3], [ChCo4] Cheeger and Colding defined a self-adjoint Laplacian operator  $\Delta$  on  $M_\infty$ . By convention  $\Delta$  is *positive*. They proved that the eigenvalues and eigenfunctions of the Laplacian  $\Delta_i$  over  $M_i^n$  converge to those on  $M_\infty$ , thereby establishing Fukaya's conjecture [Fu]. So

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if we consider  $\mathcal{RIC}$ , the completion under measured Gromov-Hausdorff convergence of the set of smooth manifolds with (0.1), it is natural to expect that quantities associated to  $\Delta$  should behave continuously.

In this paper, we will study this phenomenon in detail. Our main goal is to prove, in various cases, that the heat kernel  $H_i$  and Green's function  $G_i$  on  $M_i^n$  converge uniformly to the heat kernel  $H_\infty$  and Green's function  $G_\infty$  on  $M_\infty$ . We will make precise the definition of these convergences in Section 1. For results concerning heat kernels, a lower bound in Ricci curvature (0.1) is enough; for Green's functions, we require that  $\text{Ric}_{M_i^n} \geq 0$ ; compare with (1.16), (1.18), (1.19).

Moreover, in the noncompact case, we also study the asymptotics of the heat kernel and Green's function on a manifold  $M^n$  with  $\text{Ric}_{M^n} \geq 0$  and a Euclidean volume growth condition:

$$(0.4) \quad \text{Vol}(B_R(p)) > v_0 R^n.$$

According to [ChCo1], any tangent cone at infinity of a manifold  $M^n$  with  $\text{Ric}_{M^n} \geq 0$  and (0.4) is a metric cone  $C(X)$ . So viewed from a sufficiently large scale,  $M^n$  appears to be close to some  $C(X)$ . Combined with the appropriate convergence theorems mentioned above, at a sufficiently large scale the heat kernel and Green's function on  $M^n$  are close to those on  $C(X)$ .

On the other hand, we show that the classical analysis on cones, [Ch1], [Ch2], [ChTa1], can be generalized to  $C(X)$ . In particular, we have explicit expression of heat kernels and Green's functions on  $C(X)$ ; see (6.21), (4.23). In this way we get a unified treatment for the asymptotic formulae of these integral kernels on  $M^n$ . In particular, we get new proofs of the Colding-Minicozzi asymptotic formula for Green's functions, [CoMi1] (compare with [LiTW]), the asymptotic formulae for heat kernels of Li [Li1] and Li-Tam-Wang [LiTW].

The organization of this paper is as follows:

Section 1 reviews some background material that we need in the sequel.

In Section 2, in the *compact* case we prove,  $H_i(\cdot, \cdot, t) \rightarrow H_\infty(\cdot, \cdot, t)$  uniformly (assuming (0.1)), and  $G_i \rightarrow G_\infty$  uniformly, off the diagonal (assuming  $\text{Ric}_{M^n} \geq 0$ ). It's well known that there is an eigenfunction expansion for heat kernels, so our results follows easily from the work of Cheeger-Colding [ChCo4], [Ch3], by estimating the remainders of the eigenvalue expansions. We remark, previously in [KK1], [KK2] it was proved that a *subsequence* of  $H_i$  converges to *some* kernel on the compact metric space  $M_\infty$ .

By the Dirichlet's principle and the transplantation theorem of Cheeger [Ch3], we show in Section 3 that the uniform limit of solutions for Poisson equations is a solution for a Poisson equation, see also [Ch3], [ChCo4]. In particular, if  $\{M_i^n\}$  are noncompact,  $\text{Ric}_{M_i^n} \geq 0$ , satisfy (0.4) uniformly, then  $G_i \rightarrow G_\infty$  uniformly off the diagonal (Theorem 3.21).

We treat the heat kernels on noncompact spaces in Section 5. We assume (0.1). First, some *subsequence* of the Dirichlet heat kernels  $H_{R,i}$  on  $B_R(p_i) \subset M_i^n$  will converge to some function  $H_{R,\infty}$  on  $B_R(p_\infty) \subset M_\infty$ . However, at present it is not clear if  $H_{R,i}$  will converge. On the other hand by a generalized maximum principle, any two  $H_{R,\infty}$  (from two different subsequences) can not be too different from each other, see (5.46). Letting  $R \rightarrow \infty$ , we prove that  $H_i(\cdot, \cdot, t) \rightarrow H_\infty(\cdot, \cdot, t)$  in  $L^1$ .

In Theorem 5.59, when the noncollapsed limit  $M_\infty$  is a manifold, we prove  $H_\infty$  is the heat kernel over  $M_\infty$ , i.e. the integral kernel of the semigroup  $e^{-t\Delta}$ . For general  $M_\infty$ , the picture is not yet clear; however, it is true when  $M_\infty = C(X)$  is a noncollapsed tangent cone that is the limit of a sequence of manifolds with nonnegative Ricci curvature, see Theorem 6.1.

In Section 4 and Section 6 we study the Laplacian on  $C(X)$ . We prove that in this case, one can still separate variables. We use these to study the structure of  $G_\infty$  and  $H_\infty$  on  $C(X)$ , and the asymptotic behavior of Green's function and heat kernel on a manifold  $M^n$  with  $\text{Ric}_{M^n} \geq 0$  and (0.4).

In Section 7 we study the asymptotic behavior of the eigenvalues  $\lambda_{j,\infty}$  on a compact metric space  $M_\infty$  which is the limit of a sequence of manifolds  $\{M_i^n\}$  with (0.1). We will prove in the noncollapsed case, the Weyl asymptotic formula is true on  $M_\infty$ ; see Theorem 7.3. In the collapsed case, we get some link between the behavior of eigenvalues and  $\dim_{\text{Mink}}(M_\infty)$ , the Minkowski dimension of  $M_\infty$ .

All of the estimates in this paper are *uniform*, i.e. the constants are valid for the whole family of manifolds we are considering (for example all compact manifolds  $M^n$  with  $\text{Ric}_{M^n} \geq -(n-1)\Lambda$  and  $\text{Diam } M^n \leq D$ ).

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## 1. BACKGROUND AND NOTATION

Suppose  $(M_i^n, \text{Vol}_i) \xrightarrow{d_{GH}} (M_\infty, \mu_\infty)$  in the measured Gromov-Hausdorff sense, i.e. the sequence  $\{M_i^n\}$  converges in the Gromov-Hausdorff sense to  $M_\infty$ , and for any  $x_i \rightarrow x_\infty$ , ( $x_i \in M_i^n$ ) and  $R > 0$ , we have  $\text{Vol}_i(B_R(x_i)) \rightarrow \mu_\infty(B_R(x_\infty))$ ; here  $\mu_\infty$  is Borel regular. In fact, for any sequences of manifolds with Ricci curvature bounded from below, after possible renormalization of the measures when  $\{M_i^n\}$  is collapsing, one can always find a subsequence converges in the measured Gromov-Hausdorff sense, see [ChCo2]. In the following we usually let  $\text{Vol}$  denote the (renormalized if  $\{M_i^n\}$  is collapsing) measure on  $M_i^n$ ,  $i = 1, 2, \dots, \infty$ ; on  $M_\infty$  sometimes we also write  $\mu_\infty$  for  $\text{Vol}_\infty$ . We refer to [Ch3], [ChCo2], [Gr] for general background on measured Gromov-Hausdorff convergence.

**Definition 1.1.** *Suppose  $K_i \subset M_i^n \xrightarrow{d_{GH}} K_\infty \subset M_\infty$  in the measured Gromov-Hausdorff sense.  $f_i$  is a function on  $M_i^n$ ,  $i = 1, 2, \dots$ ;  $f_\infty$  is a continuous function on  $M_\infty$ . Assume  $\Phi_i : K_\infty \rightarrow K_i$  are  $\epsilon_i$ -Gromov-Hausdorff approximations,  $\epsilon_i \rightarrow 0$ . If  $f_i \circ \Phi_i$  converge to  $f_\infty$  uniformly, we say that  $f_i \rightarrow f_\infty$  uniformly over  $K_i \xrightarrow{d_{GH}} K_\infty$ .*

For simplicity, in the above context, we also say that  $f_i \rightarrow f_\infty$  uniformly on  $K$ ; when we write  $f_i(x) \rightarrow f_\infty(x)$ , we mean that  $f_i \rightarrow f_\infty$  uniformly and  $f_i(x_i) \rightarrow f_\infty(x_\infty)$ , where  $x_i \rightarrow x_\infty$ ,  $x_i \in M_i$ .

In many applications, the family  $\{f_i\}$  is actually equicontinuous. We remark, the Arzela-Ascoli theorem can be generalized to the case where the functions live on different spaces: when  $M_i^n \xrightarrow{d_{GH}} M_\infty$ , for any bounded, equicontinuous sequence  $\{f_i\}$

( $f_i$  is a function on  $M_i^n$ ), there is a subsequence that converges uniformly to some continuous function  $f_\infty$  on  $M_\infty$ . The proof is straightforward.

We also introduce the notion of  $L^p$  convergence ( $1 \leq p \leq \infty$ ).

**Definition 1.2.** We say  $f_i \rightarrow f_\infty$  in  $L^p$ , if for all  $\epsilon > 0$ , one can write  $f_i = \phi_i + \eta_i$  such that  $\phi_i \rightarrow \phi_\infty$  uniformly and  $\limsup_{i \rightarrow \infty} \|\eta_i\|_{L^p} \leq \epsilon$ ,  $\|\eta_\infty\|_{L^p} \leq \epsilon$ .

The following is a generalization of the Rellich-Kondrakov theorem:

**Lemma 1.3** (Rellich). Assume  $B_1(p_i) \subset M_i^n \xrightarrow{d_{GH}} B_1(p) \subset M_\infty$  in the measured Gromov-Hausdorff sense.  $u_i$  is a function on  $M_i^n$ ,  $i = 1, 2, \dots$ . Assume

$$(1.4) \quad \int_{B_1(p_i)} (u_i)^2 + |\nabla u_i|^2 \leq N.$$

Then there is a subsequence of  $\{u_i\}$  that converges in  $L^2$  over any compact subset of the open balls  $B_1$ .

The proof depends only on a *weak Poincare inequality* (use a bigger ball on the right of (0.3)) and (0.2). One can divide the ball  $B_R(p)$  ( $R < 1$ ) into small subsets and approximate  $f$  by functions that are constant over each of these small subsets, then one easily finds a convergent subsequence by standard diagonal arguments. Compare [Ch3], especially [CoMi3].

We use subscript  $i$  and write  $f_i, H_{R,i}, p_i$ , etc. to denote functions, points, etc on  $M_i^n$ . To simplify notation, when we write an equation with some function or other objects with no subscription (for example,  $f$ ), it should be understood that the equation is valid for some suitable convergent sequence of functions or other objects (for example,  $\{f_i\}$ ;  $f_i$  is defined on  $M_i^n$ ,  $i = 1, 2, \dots$ ), according to the context.

In [Ch3], Cheeger defined a Sobolev space  $H_{1,2}$  on metric-measure spaces  $(Z, \mu)$  satisfying (0.2), (0.3), and proved that Lipschitz functions are dense in  $H_{1,2}$ . Denote by  $\overset{\circ}{H}_{1,2}(\Omega)$  the closure in  $H_{1,2}$  of Lipschitz functions supported in an open set  $\Omega$ . Recall in [Ch3], one has a natural finite dimensional cotangent bundle  $T^*Z$ . We use  $du$  to denote the *differential* of  $u$ , see Section 4 of [Ch3]. One can put a norm  $|\cdot|$  on  $T^*Z$  by assigning  $|df| = g_f = \text{Lip } f$  for  $f$  Lipschitz. Here as in [Ch3],

$$(1.5) \quad \text{Lip } f(x) = \limsup_{y \rightarrow x} \frac{|f(y) - f(x)|}{d(y, x)}.$$

Note we use **Lip** $f$  to denote the Lipschitz constant of  $f$ . Clearly, on smooth manifold  $|\cdot|$  agrees with the standard norm  $|du| = |\nabla u|$ . It was proved in [Ch3] that  $|\cdot|$  is equivalent to a uniformly convex norm, in particular, an inner product.

If we have the stronger assumption  $Z = M_\infty$  with  $M_i^n \xrightarrow{d_{GH}} M_\infty$ , in [ChCo4] Cheeger and Colding proved that  $M_\infty$  is  $\mu_\infty$ -rectifiable, and, as a corollary, the norm  $|\cdot|$  actually comes from an inner product  $\langle \cdot, \cdot \rangle$ . So  $H_{1,2}$  is made into a Hilbert space:

$$(1.6) \quad \langle u, v \rangle_{H_{1,2}} = \int_{M_\infty} uv + \int_{M_\infty} \langle du, dv \rangle,$$

and for Lipschitz functions  $f$ , one has

$$(1.7) \quad \|f\|_{H_{1,2}}^2 = \|f\|_{L^2}^2 + \int_{M_\infty} |\text{Lip } f|^2.$$

Now by the standard theory of Dirichlet forms, one gets a *positive* self-adjoint Laplacian  $\Delta$  on  $M_\infty$ , see [Ch3], [ChCo4] for the details of this theory.

Recall one form of the transplantation theorem of Cheeger (for a proof, see Lemma 10.7 of [Ch3]):

**Lemma 1.8.** *Assume  $M_i^n \xrightarrow{d_{GH}} M_\infty$ .  $f_\infty$  is a Lipschitz function on  $B_R(x_\infty) \subset M_\infty$ ,  $x_i \rightarrow x_\infty$ . Then there is a sequence of Lipschitz functions  $\{f_i\}$  that converges uniformly to  $f_\infty$ , here  $f_i$  is defined on  $B_R(x_i) \subset M_i^n$ . Moreover, one can require that*

$$(1.9) \quad \limsup_{i \rightarrow \infty} \mathbf{Lip} f_i \leq \mathbf{Lip} f_\infty,$$

$$(1.10) \quad \limsup_{i \rightarrow \infty} \|\mathbf{Lip} f_i\|_{L^2} \leq \|\mathbf{Lip} f_\infty\|_{L^2}.$$

By [J], [HaKo], on length spaces satisfying (0.2), a weak Poincare inequality implies a uniform Poincare-Sobolev inequality (i.e, put  $L^{\chi p}$  norm on the left side of (0.3) for some  $\chi > 1$ ). In particular, we have a Dirichlet-Sobolev inequality for  $u \in \mathring{H}_{1,2}(B_R(p))$ . So we have

**Lemma 1.11** (Moser iteration). *If for all  $\phi$  with compact support in  $B_2(p)$ ,*

$$(1.12) \quad \int \nabla u \nabla \phi = \int cu\phi + f\phi,$$

*then for  $q > C(\kappa)$ , we have,*

$$(1.13) \quad \|u\|_{L^\infty(B_1(p))} \leq C(n, q)(1 + |c|)^{N(\tau, \kappa)}(\|u\|_{L^2(B_2(p))} + \|f\|_{L^q(B_2(p))}).$$

For proof, see chapter 4 of [Lin]. Note here we need to renormalize the measure.

Recall, on smooth manifolds we have

**Lemma 1.14** (Gradient estimate). *If  $\Delta u - cu = f$ ,  $\|u\|_{L^2} < \infty$  on  $B_{2R}(p)$  and  $f$  is a  $C^2$  function with Lipschitz constant  $\mathbf{Lip} f$ ,  $c$  is a constant, then on  $B_R(p)$  we have a gradient estimate:*

$$(1.15) \quad |\nabla u| \leq C(\|f\|_{L^\infty}, \|u\|_{L^2}, \mathbf{Lip} f)(1 + |c|)^{N(\tau, \kappa)}.$$

The proof follows a standard argument of Cheng-Yau, see [CY1], [LiY1], compare also with [Lin], [Li2], [SY]. Then use the Moser iteration to replace  $\|u\|_{L^\infty}$  with  $\|u\|_{L^2}$ .

Finally recall the Li-Yau estimates [LiY2], [SY]: If  $M^n$  satisfies (0.1), then its heat kernel  $H$  satisfies

$$(1.16) \quad H(x, y, t) \leq C(n) \text{Vol}(B_{\sqrt{t}}(x))^{-1/2} \text{Vol}(B_{\sqrt{t}}(y))^{-1/2} e^{-d^2(x,y)/5t} e^{C\Lambda t}.$$

If  $t < T$ , by volume comparison (1.16) simplifies to

$$(1.17) \quad H(x, y, t) \leq C(n, \Lambda, T) \text{Vol}(B_{\sqrt{t}}(x))^{-1} e^{-d^2(x,y)/5t} e^{C\Lambda t} t^{-C(n)} e^{C(n, \Lambda, T)d(x,y)}.$$

If we assume  $\text{Ric}_{M^n} \geq 0$ , then

$$(1.18) \quad \frac{C^{-1}(n)}{\text{Vol}(B_{\sqrt{t}}(x))} e^{-d^2(x,y)/3t} \leq H(x, y, t) \leq \frac{C(n)}{\text{Vol}(B_{\sqrt{t}}(x))} e^{-d^2(x,y)/5t}.$$

Assume  $M^n$  is noncompact, write  $G$  for the minimal positive Green's function on  $M^n$ . If  $n \geq 3$ ,  $M^n$  satisfies (0.4) and  $\text{Ric}_{M^n} \geq 0$ , then  $G$  exists ([LiY2], [LiT1], [LiT2], [SY]) and satisfies

$$(1.19) \quad C^{-1}(v_0)d(x, y)^{2-n} \leq G(x, y) \leq C(v_0)d(x, y)^{2-n}.$$

For proofs of the above estimates, see [LiY2], [SY].

## 2. THE COMPACT CASE

In this section, we assume  $M_\infty$  is compact. First we study the heat kernel  $H$ .

It's well known that if  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$ , there is a lower bound for the  $k$ th eigenvalue  $\lambda_{k,i}$  of the Laplacian  $\Delta_i$  over  $M_i^n$ :

$$(2.1) \quad \lambda_{k,i} \geq C(n, \Lambda, D)k^{\frac{2}{n}},$$

here  $\text{Diam } M_i^n \leq D$ . The proof uses only (0.2) and (0.3); see [Gr] and Theorem 4.8 of [Ch3]. Hence over  $M_i^n$  we have

$$(2.2) \quad H_i(x, y, t) = \sum_{j=0}^{\infty} e^{-\lambda_{j,i}t} \phi_{j,i}(x) \phi_{j,i}(y),$$

here  $\phi_{j,i}$  is the eigenfunction of the  $j$ th eigenvalue  $\lambda_{j,i}$ . By the Cheeger-Colding spectral convergence theorem [ChCo4], for each  $j$ ,  $\lambda_{j,i} \rightarrow \lambda_{j,\infty}$ , and  $\phi_{j,i} \rightarrow \phi_{j,\infty}$  uniformly when  $i \rightarrow \infty$ , here  $\lambda_{j,\infty}$  and  $\phi_{j,\infty}$  are the  $j$ -th eigenvalue and (renormalized) eigenfunction of  $\Delta$  on  $M_\infty$ . So (2.1) is also true for  $\lambda_{j,\infty}$ . Moreover,

$$(2.3) \quad \|\phi_{j,i}\|_{L^\infty} \leq C_1(n, \Lambda, D)(1 + \lambda_{j,i})^{C(n)} \|\phi_{j,i}\|_{L^2},$$

$$(2.4) \quad \|\nabla \phi_{j,i}\|_{L^\infty} \leq C_0(n, \Lambda, D)(1 + \lambda_{j,i})^{C(n)} \|\phi_{j,i}\|_{L^\infty}.$$

These are implied by (the proof of) the Moser iteration and the gradient estimate, see [LiY1]. By [ChCo4], (2.3), (2.4) can pass to  $M_\infty$  (on  $M_\infty$  (2.4) becomes an estimate for  $\mathbf{Lip} \phi_{j,\infty}$ ). So it makes sense to write

$$(2.5) \quad H_\infty(x, y, t) = \sum_{j=0}^{\infty} e^{-\lambda_{j,\infty}t} \phi_{j,\infty}(x) \phi_{j,\infty}(y).$$

By (2.1),  $H_\infty$  is the heat kernel over  $M_\infty$ .

Apply (2.3), (2.4) to  $\sum_{j=k}^{\infty} e^{-\lambda_{j,i}t} \phi_{j,i}(x) \phi_{j,i}(y)$ , the tail of (2.2), one easily get

**Theorem 2.6.** *Assume  $M_i^n \xrightarrow{d_{GH}} M_\infty$ ,  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$ . When  $t > 0$  fixed,  $H_i$  converges to the heat kernel  $H_\infty$  over  $M_\infty$  uniformly.  $H_\infty$  is continuous in  $t, x, y$ ; when  $t$  fixed, it is Lipschitz in  $x, y$ .*

**Corollary 2.7.** *For  $H_\infty$  on  $M_\infty$ , the Li-Yau estimate (1.18) is true if  $\text{Ric}_{M_i^n} \geq 0$ ; if  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$  then (1.16) is true.*

Next we study the Green's functions. Assume  $\text{Ric}_{M_i^n} \geq 0$ . Recall,

$$(2.8) \quad G_i(x, y) = \int_0^\infty h_i(x, y, t) dt,$$

here  $h_i(x, y, t) = H_i(x, y, t) - \phi_{0,i}(x)\phi_{0,i}(y) = H_i(x, y, t) - 1$ . Note, since the sequence  $\{M_i^n\}$  might collapse, we have to renormalize the measures and eigenvalues such that  $\text{Vol}(M_i^n) = 1$  and  $\{\phi_{j,i}\}_j$  is orthonormal, e.g.  $\lambda_{0,i} = 0$ ,  $\phi_{0,i} = 1$ . So

$$(2.9) \quad G_i(x, y) = \int_0^\epsilon h_i(x, y, t)dt + \sum_{j=1}^\infty e^{-\epsilon\lambda_{j,i}/2} \int_{\frac{\epsilon}{2}}^\infty e^{-\lambda_{j,i}t} \phi_{j,i}(x)\phi_{j,i}(y)dt.$$

When  $x$  is fixed, by (2.3),

$$(2.10) \quad \sum_{j=k}^\infty \|e^{-\epsilon\lambda_{j,i}/2} \int_{\frac{\epsilon}{2}}^\infty e^{-\lambda_{j,i}t} \phi_{j,i}(x)\phi_{j,i}(y)dt\|_{L^\infty} \leq \sum_{j=k}^\infty \frac{e^{-\epsilon\lambda_{j,i}}}{\lambda_{j,i}} C_1(1 + \lambda_{j,i})^{C_3}.$$

This goes to 0 uniformly in  $i$  as  $k \rightarrow \infty$ , by (2.1). On the other hand, clearly (1.18) holds after renormalization; when  $\text{Ric} \geq 0$ ,  $R < 1/8$  we have the (rescaled) volume bound

$$(2.11) \quad C(n, D)R^n \leq \text{Vol}(B_R(x)) \leq \sqrt{R}C(n) \text{Vol}(B_{\sqrt{R}}(x)) \leq C'(n)R.$$

So when  $d(x, y) \geq \delta > 0$ , by (1.18),

$$(2.12) \quad \int_0^\epsilon |h_i(x, y, t)|dt \leq \epsilon + \int_0^\epsilon |H_i(x, y, t)|dt \leq \epsilon + C \int_0^\epsilon t^{-\frac{n}{2}} e^{-\frac{\delta^2}{5t}} dt.$$

So in particular, by choosing  $\epsilon$  small, we get a function  $G_\infty(x, y)$  on  $M_\infty$ , such that  $G_i \rightarrow G_\infty$  in  $L^\infty$  on compact subsets, off the diagonal.

Finally, we want to check  $G_\infty$  is the Green's function over  $M_\infty$ . We now establish an  $L^1$  bound for  $G(x, y)$  over the ball  $B_R(x)$ . Note

$$(2.13) \quad \int_{B_R(x)} |G(x, y)|dy \leq \int_{B_R(x)} \int_0^1 |h(x, y, t)|dtdy + \int_{B_R(x)} \int_1^\infty |h(x, y, t)|dtdy.$$

Since  $\|\phi(y)\|_{L^2} = 1$ , by (2.1), (2.3), (2.11) and the Schwartz inequality,

$$(2.14) \quad \begin{aligned} \int_{B_R(x)} \int_1^\infty |h(x, y, t)|dtdy &\leq \sum_{j=1}^\infty \int_{B_R(x)} e^{-\lambda_j/2} |\phi_j(x)| \int_{\frac{1}{2}}^\infty e^{-\lambda_j t} |\phi_j(y)|dtdy \\ &\leq C(n) \sum_{j=1}^\infty e^{-\lambda_j/2} (1 + \lambda_j)^{C_2(n)} \sqrt{R} \int_{\frac{1}{2}}^\infty e^{-\lambda_j t} dt \leq C'(n) \sqrt{R}. \end{aligned}$$

Now we focus on the first term on the right hand side of (2.13). Since  $H - h = 1$ , and we have (2.11), it's enough to estimate the integral of  $H$ . Put  $R < 1/8$ , by (1.18), (2.11) we have

$$(2.15) \quad \begin{aligned} \int_0^1 \int_{B_R(x)} H(x, y, t)dtdy &\leq \int_0^1 \int_{B_R(x)} C(n) \text{Vol}^{-1}(B_{\sqrt{t}}(x)) e^{-\frac{d(x,y)^2}{5t}} dtdy \\ &\leq \left( \int_0^R \int_{B_R(x)} C(n) \text{Vol}^{-1}(B_{\sqrt{t}}(x)) e^{-\frac{d(x,y)^2}{5t}} dtdy \right) + C'(n) \sqrt{R}. \end{aligned}$$

Next,

$$(2.16) \quad \int_0^R \frac{\int_{B_R(x)} e^{-d(x,y)^2/5t} dy}{\text{Vol}(B_{\sqrt{t}}(x))} dt = \int_0^R \frac{\int_0^R e^{-r^2/5t} A(r) r^{n-1} dr}{\int_0^{\sqrt{t}} A(r) r^{n-1} dr} dt,$$

here  $A(r)r^{n-1}$  is the surface area element of  $\partial B_r(x)$ . Since  $\text{Ric}_{M_i^n} \geq 0$ ,  $A(r)$  is non-increasing. The right hand side of (2.16) can be bounded by

$$(2.17) \quad \int_0^R \left( \sum_{s=0}^{\lfloor R/\sqrt{t} \rfloor + 1} \int_{s\sqrt{t}}^{(s+1)\sqrt{t}} e^{-r^2/5t} A(r)r^{n-1} dr / \int_0^{\sqrt{t}} A(r)r^{n-1} dr \right) dt \\ \leq C_1(n)R + \int_0^R \sum_{s=1}^{\lfloor R/\sqrt{t} \rfloor + 1} e^{-s^2/5} ((s+1)^n - s^n) dt \leq C'(n)R.$$

So combine (2.14), (2.15) and (2.11) we get

$$(2.18) \quad \int_{B_R(x)} |G(x, y)| dy \leq C'(n)\sqrt{R}.$$

Since  $G_i \rightarrow G_\infty$  uniformly off the diagonal, use the Cheeger-Colding theorem on the convergence of eigenfunctions [ChCo4] and (2.3), (2.4), we get, for all  $x$ ,

$$(2.19) \quad \phi_{j,\infty}(x) = \lim_{i \rightarrow \infty} \phi_{j,i}(x) = \lim_{i \rightarrow \infty} \int_{M_i} G_i(x, y) \lambda_{j,i} \phi_{j,i}(y) dy \\ = \int_{M_\infty} G_\infty(x, y) \lambda_{j,\infty} \phi_{j,\infty}(y) dy.$$

So  $G_\infty$  is the Green's function over  $M_\infty$ . Moreover, by (2.10), (2.12) and Lemma 1.14,  $G_\infty$  is Lipschitz continuous off the diagonal. It is harmonic off the diagonal by Lemma 3.17. So we have proved

**Theorem 2.20.** *Assume  $M_i^n \xrightarrow{d_{GH}} M_\infty$ ,  $\text{Ric}_{M_i^n} \geq 0$ . Then the Green's function  $G_\infty$  on  $M_\infty$  exists. On any compact subsets  $K$  off the diagonal,  $G_\infty$  is Lipschitz and harmonic,  $G_i \rightarrow G_\infty$  uniformly on  $K$ .*

### 3. THE GREEN'S FUNCTIONS ON NONCOMPACT SPACES

Recall how on a manifold, one solves the the Poisson equation,

$$(3.1) \quad \Delta u_R = f, \quad u_R|_{\partial B_R(p)} = h,$$

for Lipschitz functions  $f, h$  on the closed ball  $B_R(p)$ . By the Dirichlet's principle,  $u_R$  is the unique minimizer of the functional

$$(3.2) \quad I(u) = \int_{B_R(p)} (|du|^2 - fu).$$

within the space  $\mathcal{E} = h + \mathring{H}_{1,2}(B_R(p))$ , note  $\Delta$  is positive by convention.

Assume  $M_i^n \xrightarrow{d_{GH}} M_\infty$  in the measured Gromov-Hausdorff sense,  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$ . Recall, by [Ch3], [ChCo4],  $\Delta$  is linear on  $M_\infty$ . So the above variational method is valid also on  $M_\infty$ .

**Lemma 3.3** (Lower semicontinuity of energy). *Suppose  $u_i, f_i$  are  $C^2$  functions over  $M_i^n$ ,  $\Delta u_i = f_i$ ,  $u_i \rightarrow u_\infty$ ,  $f_i \rightarrow f_\infty$  uniformly over the sequence of converging balls  $B_{2R}(p_i) \rightarrow B_{2R}(p_\infty)$ , and there is a uniform gradient estimate for  $u_i$  and  $f_i$ :*

$$(3.4) \quad |\nabla u_i|, |\nabla f_i| < L.$$

Then we have

$$(3.5) \quad I(u_\infty) \leq \liminf_{i \rightarrow \infty} I(u_i).$$

*Proof.* As in [ChCo1], we can get an integral bound for the Hessian of  $f_i$  on the ball  $B_1(p_i)$ : recall the Bochner formula

$$(3.6) \quad \frac{1}{2} \Delta(|\nabla f_i|^2) = |\text{Hess}_{f_i}|^2 + \langle \nabla \Delta f_i, \nabla f_i \rangle + \text{Ric}(\nabla f_i, \nabla f_i).$$

Multiply by a cut-off function  $\phi$  with  $\text{supp} \phi \subset B_r \subset B_1(q_i)$ ,  $\phi|_{B_{r/2}} = 1$ ,  $|\nabla \phi| \leq c(n, r)$ ,  $|\Delta \phi| \leq c(n, r)$ ; see Theorem 6.33 of [ChCo1]. Since  $f_i$  is harmonic,

$$(3.7) \quad \frac{1}{2} \phi \Delta(|\nabla f_i|^2) = \phi |\text{Hess}_{f_i}|^2 + \phi \text{Ric}(\nabla f_i, \nabla f_i).$$

Integrate by parts,

$$(3.8) \quad \int_{B_r} \frac{1}{2} (|\nabla f_i|^2) \Delta \phi = \int_{B_r} \phi |\text{Hess}_{f_i}|^2 + \int_{B_r} \phi \text{Ric}(\nabla f_i, \nabla f_i).$$

By assumption, there is a definite lower bound for the last term in (3.8). Note  $|\Delta \phi|$  is uniformly bounded by construction, we have a uniform upper bound for  $\int_{B_r} \phi |\text{Hess}_{f_i}|^2$ . So by Lemma 1.3 we can assume some subsequence of  $|\nabla f_i|$  converges to a function  $\Gamma$  on  $B_R(p_\infty) \subset M_\infty$  in  $L^2$ . Assume,  $x \in \mathcal{R}_k \subset M_\infty$  for some  $k$  (all tangent cone at  $x$  is  $\mathbf{R}^k$ ), there is some subset  $A(x) \subset M_\infty$  such that  $\Gamma$  is continuous on  $A(x)$ ,  $x \in A(x)$  is a density point of  $A(x)$ . By Luzin's theorem and the results in [ChCo2], these properties hold for almost all  $x \in M_\infty$ . For such  $x$ , we prove

$$(3.9) \quad |\text{Lip } f_\infty(x)| \leq \Gamma(x).$$

Clearly, (3.9) implies our lemma.

To prove (3.9), it's enough to prove, for all  $\psi > 0$ , if  $l = d(x, y)$  is sufficiently small, then

$$(3.10) \quad |f_\infty(x) - f_\infty(y)| \leq d(y, x)(\Gamma(x) + 6\psi).$$

By the gradient estimate of  $f_i$  (so of  $f_\infty$ ), if (3.10) is not true for some  $y_0$ , then for all  $y \in B_{l\psi/L}(y_0)$ ,

$$(3.11) \quad |f_\infty(x) - f_\infty(y)| > d(y, x)(\Gamma(x) + 5\psi).$$

Pick  $x_i, y_{0,i} \in M_i^n$ ,  $x_i \rightarrow x$ ,  $y_{0,i} \rightarrow y_0$ ,  $d(x_i, y_{0,i}) = l$ . Then for  $i$  big enough, for all  $y_i \in B_{l\psi/L}(y_{0,i})$  and all minimal geodesic  $\gamma_i$  connecting  $x_i$  and  $y_i$ ,

$$(3.12) \quad \int_{\gamma_i} |\nabla f_i| \geq d(x_i, y_i)(\Gamma(x) + 4\psi).$$

First of all, since  $|\nabla f_i|$  is uniformly bounded by  $L$ , a simple computation shows along every  $\gamma_i$  we must have

$$(3.13) \quad |\nabla f_i| > \Gamma(x) + 2\psi,$$

on a subset of  $\gamma_i$  which has 1-Hausdorff measure at least  $2\psi l / (L - \Gamma(x))$ . Put

$$(3.14) \quad T_i = \{v \in T_{x_i} | v = \gamma'(0) \text{ for some minimal geodesic } \gamma \text{ from } x_i \text{ to } y_i \in B_{l\psi/L}(y_{0,i})\}.$$

We must have

$$(3.15) \quad H^{n-1}(T_i) > C(n, L, \psi)H^{n-1}(\partial B_1(0)),$$

where  $H^{n-1}$  is the  $(n-1)$ -Hausdorff measure over the unit sphere  $\partial B_1(0)$  in the tangent space  $T_{x_i}$ . Combine this with (3.13), when  $M_\infty$  is *noncollapsed*, if  $l$  is small enough, by the proof of the Bishop-Gromov inequality, for sufficiently big  $i$ ,

$$(3.16) \quad \frac{\text{Vol}(\{z_i \in B_l(x_i) \mid |\nabla f_i(z_i)| > \Gamma(x) + 2\psi\})}{\text{Vol}(B_l(x_i))} \geq C(x, n, L, \psi) > 0.$$

Now  $|\nabla f_i|$  converge to  $\Gamma$  in  $L^2$ , so (3.16) is also true if we substitute  $|\nabla f_i|$  in (3.16) by  $\Gamma$ ,  $x_i$  by  $x$ . We get a contradiction to the choice of  $x$ .

The proof is the same when  $M_\infty$  is *collapsed*. We just use the *segment inequality* ([ChCo1], [ChCo4]) to get (3.16) from (3.11), (3.12) and (3.13).  $\square$

**Lemma 3.17.** *Let  $u_\infty, f_\infty$  be as in the previous lemma. Then*

$$(3.18) \quad \Delta u_\infty = f_\infty.$$

*Proof.* Assume this is not true over a ball  $B_\lambda(p^*) \subset \subset B_1(0)$ . By solving the Dirichlet problem on  $B_\lambda(p^*)$  we can find  $v_\infty$  with the same boundary value as  $u_\infty$  over  $\partial B_\lambda$ , but with smaller energy, say

$$(3.19) \quad I(v_\infty) = \int_{B_\lambda(p^*)} |dv_\infty|^2 - f_\infty v_\infty < \int_{B_\lambda(p^*)} |du_\infty|^2 - f_\infty u_\infty - 2\Psi.$$

By obvious density properties, we can change  $v_\infty$  slightly so that  $v_\infty$  agrees with  $u_\infty$  on a neighborhood of  $\partial B_\lambda(p^*)$ . By Lemma 3.3, for  $i$  big enough,

$$(3.20) \quad I(v_\infty) \leq I(u_i) - \Psi.$$

So by (the proof of) Lemma 1.8 (see Section 10 of [Ch3]), for  $i$  big enough we can find a function  $v_i$  with the same boundary value on  $\partial B_i$  as  $u_i$  but with smaller energy  $I$ . That contradicts the fact that  $\Delta u_i = f_i$ . In view of (3.2) and Lemma 1.8.  $\square$

The solution of (3.1) is unique on  $M_\infty$  because the maximum principle holds, see Section 7 of [Ch3].

We now study the Green's functions. Assume  $(M_i^n, p_i, \text{Vol}_i) \xrightarrow{d_{GH}} (M_\infty, p, \mu_\infty)$  in the pointed measured Gromov-Hausdorff sense ([Gr], [ChCo2]), where  $\text{Ric}_{M_i^n} \geq 0$ ,  $M_i^n$  is complete, noncompact,  $n \geq 3$ .

**Theorem 3.21.** *Assume,  $M_i^n$  also satisfies  $\text{Vol}(B_R(p_i)) > v_0 R^n$ . Then on  $M_\infty$  there is a Green's function  $G_\infty$ ,  $G_i \rightarrow G_\infty$  uniformly on any compact subsets of  $M \times M$  that does not intersect with the diagonal.*

*Proof.* Since  $n \geq 3$ , the Euclidean volume growth condition (0.4) implies that the minimal positive Green's function  $G_i$  exists on  $M_i^n$  ( $\Delta$  is *positive*). Moreover,  $G_i$  satisfies the Li-Yau estimate (1.19). So by the Cheng-Yau gradient estimate and the Arzela-Ascoli theorem, for any fixed  $x$ , for some subsequence (still denoted by  $G_i$ ), we have

$$(3.22) \quad G_i(x, y) \rightarrow G_\infty(x, y),$$

uniformly over any compact set in  $M \setminus \{x\}$ . Clearly,  $G_\infty$  satisfies (1.19). We will show  $G_\infty$  is in fact well defined, and  $G_i \rightarrow G_\infty$  as stated in the above theorem.

Assume  $f_\infty$  is any Lipschitz function supported in  $B_K(p_\infty) \subset M_\infty$ ,  $\mathbf{Lip} f_\infty \leq L$ . By Lemma 1.8 and approximation, there is a sequence of  $C^2$  functions  $\{f_i\}$  with  $f_i \rightarrow f_\infty$  uniformly,  $\mathbf{Lip} f_i \leq 2L$ ,  $\text{supp} f_i \subset B_{2K}(p_i) \subset M_i^n$ ,  $i = 1, 2, \dots, \infty$ .

Recall on each manifold  $M_i^n$  with maximal volume growth condition, the function,

$$(3.23) \quad u_i(x) = \int_{M_i^n} G_i(x, y) f_i(y) dy,$$

solves the Poisson equation

$$(3.24) \quad \Delta u_i = f_i, \quad \lim_{x \rightarrow \infty} u_i(x) = 0.$$

Now by the Li-Yau estimate (1.19) and the Euclidean volume growth condition (0.4),  $G_i$  is locally integrable, so  $u_i$  is uniformly bounded. The gradient estimate Lemma 1.14 shows that  $u_i$  are uniformly Lipschitz:

$$(3.25) \quad \mathbf{Lip} u_i \leq C(L, K, n).$$

So we can find a subsequence of  $\{u_i\}$  that converges to some Lipschitz function  $u_\infty$  on  $M_\infty$ . Note that by the Li-Yau estimate, (1.19),

$$(3.26) \quad |u_i(x)| \leq C'(L, K, n) d(x, p_i)^{2-n}, \quad (i = 1, 2, \dots, \infty).$$

So by Lemma 3.17,  $\Delta u_\infty = f_\infty$  on  $M_\infty$ . Using the fact Laplacian is linear, by (3.26) and the maximum principle (Section 7 of [Ch3]), it is clear that  $u_\infty$  is well defined and  $u_i \rightarrow u_\infty$  uniformly.

Notice, by (1.19),

$$(3.27) \quad u_\infty(x) = \int_{M_\infty} G_\infty(x, y) f_\infty(y) dy.$$

Since we can choose arbitrary  $K$ ,  $f_\infty$ , clearly  $G_\infty$  is also well defined,  $G_i \rightarrow G_\infty$  uniformly, off the diagonal. By (3.27) and Lemma 3.17,  $G_\infty$  can be interpreted as the minimal positive Green's function on  $M_\infty$ .  $\square$

#### 4. SEPARATION OF VARIABLES ON TANGENT CONES

Assume  $M_i^n$  is complete noncompact,  $\text{Ric}_{M_i^n} \geq 0$  and satisfies (0.4) uniformly,  $M_i^n \xrightarrow{d_{GH}} M_\infty$ . Recall that by [ChCo1], [ChCo2], every tangent cone of  $M_\infty$  is a metric cone. We denote such a cone by  $C(X) = \mathbf{R}_+ \times_r X$ , here  $(X, dx^2)$  is a compact length space with  $\text{Diam} X \leq \pi$ , [ChCo1]. The metric on  $C(X)$  is

$$(4.1) \quad d\rho^2 = dr^2 + r^2 dx^2.$$

Here we write  $r$  for the distance from the pole  $p_\infty = (0, X)$ .

The measure  $\mu_\infty$  on  $C(X)$  is just the  $n$ -Hausdorff measure, [ChCo2]. Since we can rescale  $C(X)$ ,  $\mu_\infty$  induced a natural measure  $\mu_X$  on  $X$  that obviously satisfies a doubling condition (0.2) (with some different  $\kappa$ ). Moreover,  $X$  satisfies the rectifiability properties as stated in Section 5 of [ChCo4].

Also recall from [Ch3], for  $f, g \in H_{1,2}$ ,

$$(4.2) \quad d(fg) = f \cdot dg + g \cdot df.$$

Moreover, from [ChCo4] and [Ch3], if  $f$  is a function depending only on  $r$  and  $g$  is a function independent of  $r$ , then by the polar identity, one gets  $\langle df, dg \rangle = 0$ .

**Lemma 4.3** (Weak Poincare inequality). *For  $B_R \subset X$ ,  $3R < \frac{1}{5}$ ,  $f \in H_{1,2}(X)$ ,*

$$(4.4) \quad \int_{B_R(x)} |f - f_{x,R}|^2 \leq \tau_X R^2 \int_{B_{3R}(x)} |df|^2.$$

*Proof.* Define, for  $x \in X$ ,

$$(4.5) \quad \text{Box}((1, x), a, b) = \{(t, y) \in C(X) \mid |t - 1| < a, d_X(x, y) < b\}.$$

Put

$$(4.6) \quad \text{Box}_1 = \text{Box}((1, x), R, R), \quad \text{Box}_2 = \text{Box}((1, x), 3R, 3R) \subset C(X)$$

So  $\text{Box}_1 \subset B_{2R}((1, x)) \subset \text{Box}_2$ . We extend  $f$  to be a  $H_{1,2}$  function independent of  $r$  on  $C(X)$ . Assume  $f_{\text{Ball}}$  is the average of  $f$  on the ball  $B_{2R}((1, x)) \subset C(X)$ ,

$$(4.7) \quad \begin{aligned} \int_{B_R(x)} |f - f_{x,R}|^2 &= C(n)R^{-1} \int_{\text{Box}_1} |f - f_{x,R}|^2 \leq C(n)R^{-1} \int_{\text{Box}_1} |f - f_{\text{Ball}}|^2 \\ &\leq C(n)R^{-1} \int_{B_{2R}((1, x))} |f - f_{\text{Ball}}|^2 \leq C(n)\tau R \int_{B_{2R}((1, x))} |df|^2 \\ &\leq C(n)\tau R \int_{\text{Box}_2} |df|^2 = \tau_X R^2 \int_{B_{3R}(x)} |df|^2. \end{aligned}$$

The first and last identity come from the Fubini theorem. Note  $f_{x,R}$  is also the average of  $f$  over  $\text{Box}_1$ , and we used the Poincare inequality on  $C(X)$  in the middle inequality.  $\square$

We remark, a weak Poincare inequality is already enough for many purposes. Since  $X$  is a length space, by [HaKo] one has (0.3) on  $X$ . As in [Ch3], [ChCo4], we define a *positive* operator  $\Delta_X$  on  $X$ . Note by (0.2), (0.3) the compact embedding lemma 1.3 is true on  $X$ . So by the standard elliptic theory, on  $X$  we have a basis  $\{\phi_j\}_{j=0}^\infty$  for  $L^2(X)$  and a sequence  $\mu_j \rightarrow \infty$  such that  $\Delta_X \phi_j = \mu_j \phi_j$ , compare [Ch3], [ChCo4]. Moreover, one can do Moser iteration on  $X$ , so  $\phi_i$  is Hölder continuous; see [Lin], [GT]. These have applications in Section 6.

Next we show, even the cross section  $X$  may not be a manifold, there is still a separation of variables formula for  $\Delta$  on  $C(X)$ . See [Ch1] for the classical case.

Recall that  $\langle \cdot, \cdot \rangle$  is the inner product on  $T^*M_\infty$  as in [Ch3], [ChCo4],

**Lemma 4.8.**

$$(4.9) \quad \Delta(fg) = f\Delta g + g\Delta f - 2 \langle df, dg \rangle.$$

*Proof.* Since  $d(fg) = f \cdot dg + g \cdot df$ , for any Lipschitz (or  $H_{1,2}$ ) function  $\phi$  with compact support, we have (recall  $\Delta$  is *positive*)

$$(4.10) \quad \int \langle df, g \cdot d\phi + \phi \cdot dg \rangle - \int g\phi\Delta f = 0.$$

Exchange the role of  $f$  and  $g$ , we get

$$(4.11) \quad \int \langle d(fg), d\phi \rangle - \int \phi(f\Delta g + g\Delta f - 2 \langle df, dg \rangle) = 0.$$

□

Similarly, by  $d(f \circ g) = f'(g)dg$ , we get

$$(4.12) \quad \Delta f \circ g = -f''(g)|dg|^2 + f'(g)\Delta g.$$

**Lemma 4.13.** *On  $C(X)$ ,  $r^{2-n}$  is harmonic away from the pole.*

*Proof.* By the results in Section 4 of [ChCo1],  $r^{2-n}$  is the uniform limit of a sequence of harmonic functions  $\mathcal{G}$ . So by the proof of Lemma 3.17,  $r^{2-n}$  is harmonic. □

By the maximum principle on  $X$  (Section 7 in [Ch3]), we have

**Lemma 4.14.** *If  $X$  is compact, and  $\Delta_X f = 0$ , then  $f$  is a constant.*

**Theorem 4.15.** *Assume  $u$  lies in the ring generated by functions of the form  $u = fg$  where  $f$  depends only on  $r$  and  $g$  depends only on  $x$ . Then on  $C(X) \setminus \{p_\infty\}$ ,*

$$(4.16) \quad \Delta u = -\frac{\partial^2 u}{\partial r^2} - \frac{n-1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \Delta_X u.$$

*Proof.* Compare [ChTa1]. By (4.12) and Lemma 4.13, on the cone  $C(X)$  we have

$$(4.17) \quad \Delta f(r) = -f''(r) - f' \frac{n-1}{r}.$$

Next we apply Lemma 4.8, recall  $\langle df, dg \rangle = 0$ . We pick a test function  $\phi$  of the form  $\phi = a(r)b(x)$ . By scaling we see,  $\Delta g(R, x) = R^{-2} \Delta g(1, x)$ . Assume  $a$  is supported over the interval  $[\alpha, \beta]$ ,

$$(4.18) \quad \int_{C(X)} \langle dg, d\phi \rangle = \int_\alpha^\beta (t^{1-n} \int_X t^{-2} \langle dg, a(t)db \rangle dx) dt,$$

here in the second integral we view  $g$  and  $b$  as functions on the cross section  $X = (1, X)$ . So we compute

$$(4.19) \quad \int_{C(X)} \langle dg, d\phi \rangle = \int_\alpha^\beta (t^{-n-1} \int_X a(t)b(x)\Delta_X g dx) dt. = \int \phi r^{-2} \Delta_X g.$$

Since we can choose arbitrary  $a, b$ , and  $\Delta g(R, x) = R^{-2} \Delta g(1, x)$ , we get

$$(4.20) \quad \Delta g(R, x) = R^{-2} \Delta_X g.$$

This suffices to complete the proof. □

Using transformation  $D_R : (r, x) \mapsto (Rr, x)$ , we deduce from the existence and uniqueness of  $G_\infty$  that

$$(4.21) \quad G_\infty(D_R x, D_R y) = R^{2-n} G_\infty(x, y).$$

So  $G_\infty(p_\infty, x) = d(p_\infty, x)^{2-n} g(x)$  for some Lipschitz function  $g$ . By (4.16) and Lemmas 4.8, 4.14,  $g = C$  is a constant.

**Corollary 4.22.**

$$(4.23) \quad G_\infty(p_\infty, x) = (n-2)^{-1} \mu_X(X)^{-1} d^{2-n}(p_\infty, x).$$

*Proof.* We know  $G_\infty(p_\infty, x) = Cd^{2-n}(p_\infty, x)$ . We construct a test function  $\phi = \phi(r)$  such that  $\phi$  is a smooth function of  $r$ ,  $\phi = 1$  for  $r$  small and  $\phi = 0$  for  $r \geq 1$ . So

$$(4.24) \quad \begin{aligned} 1 &= \int G(p_\infty, y) \Delta \phi(y) = \int_0^1 \left(-\phi'' - \frac{n-1}{r} \phi'\right) Cr^{2-n} r^{n-1} \mu_X(X) dr \\ &= -C\mu_X(X) \int_0^1 (n-2)\phi' dr = (n-2)C\mu_X(X). \end{aligned}$$

□

**Corollary 4.25.** *Assume  $\Delta u = f$  on  $B_R(p_\infty) \setminus \{p_\infty\}$ ,  $f \in L^\infty$ , and*

$$(4.26) \quad \lim_{x \rightarrow p} |u(x)| d(x, p_\infty)^{n-2} = 0.$$

*Then  $\Delta u = f$  on  $B_R(p_\infty)$ .*

*Proof.* By the De Giorgi-Nash-Moser theorem,  $u$  is bounded and Hölder continuous. In our case,  $G_\infty(p_\infty, x) = Cd(p_\infty, x)^{2-n}$ , so the proof goes exactly like the  $\mathbf{R}^n$  case (where the maximum principle is used). For details see [Lin]. □

Relation (4.23) implies the Colding-Minicozzi asymptotic formula, [CoMi1], compare [LiTW]. In fact, we rescale the manifold  $M^n$  to get a sequence of manifolds that converges to  $C(X)$ , a tangent cone at infinity, see [ChCo2]. By Theorem 3.21, the new (rescaled) Green's functions converge to the Green's function on  $C(X)$ .

**Theorem 4.27** (Colding-Minicozzi). *On a noncompact manifold  $M^n$  with  $\text{Ric}_{M^n} \geq 0$  and (0.4) we have*

$$(4.28) \quad \lim_{d(x,p) \rightarrow \infty} d(x,p)^{n-2} G(p,x) = (n-2)^{-1} \left( n \lim_{R \rightarrow \infty} R^{-n} \text{Vol}(B_R(p)) \right)^{-1}.$$

Note the tangent cones may not be unique; in collapsing case, a tangent cone might not be a metric cone, [ChCo2], [Per].

## 5. HEAT KERNELS ON NONCOMPACT SPACES

We assume in this section, all the manifolds  $M^n$  are noncompact, satisfying (0.1). On  $M^n$ , write  $H(x, y, t)$  for the heat kernel; we denote by  $H_R(x, y, t)$  the Dirichlet heat kernel on the metric ball  $B_R(p)$ , put  $H_R = 0$  outside  $B_R(p)$ .

One technical issue is, the boundary  $\partial B_R(p) = d^{-1}(R)$  may not be smooth, here  $d = d(p, \cdot)$ . However, we can approximate  $d$  by a Morse function  $d_\epsilon$ , see [Hir], and (assuming)  $R$  is not a critical value, etc. So in the sequel we always assume the boundary are smooth.

**Lemma 5.1.** *Assume  $\text{Ric}_{M^n} \geq -(n-1)\Lambda$ . Then there is a function  $\epsilon(t, \Lambda, R)$  with  $\lim_{R \rightarrow \infty} \epsilon(t, \Lambda, R) = 0$  for  $t > 0$ , and*

$$(5.2) \quad \int_{M-B_R(x)} H(x, y, t) dy \leq \epsilon(t, \Lambda, R).$$

*Proof.* By the Bishop-Gromov inequality, it's easy to see

$$(5.3) \quad \text{Vol}(B_{\sqrt{t}}(x)) \leq C_1(n, \Lambda, t) e^{C_2(n, \Lambda, t)d(x, y)} \text{Vol}(B_{\sqrt{t}}(y)).$$

Put  $s_\Lambda(r) = (1/\sqrt{\Lambda}) \sinh \sqrt{\Lambda}r$ . We now use the Li-Yau estimate (1.16):

$$(5.4) \quad \begin{aligned} & \int_{M-B_R(x)} H(x, y, t) dy \\ & \leq C'(n, \Lambda, t) \int_{M-B_R(x)} \text{Vol}^{-1}(B_{\sqrt{t}}(x)) e^{-d(x, y)^2/5t} e^{C_2(n, \Lambda, t)d(x, y)} \\ & = C'(n, \Lambda, t) \int_R^\infty e^{-r^2/5t} e^{C_2 r} A(r) s_\Lambda^{n-1}(r) dr / \int_0^{\sqrt{t}} A(r) s_\Lambda^{n-1}(r) dr \\ & \leq C' \int_R^\infty e^{-r^2/5t} e^{C_2 r} s_\Lambda^{n-1}(r) dr / \int_0^{\sqrt{t}} s_\Lambda^{n-1}(r) dr = \epsilon(t, \Lambda, R). \end{aligned}$$

Here  $A(r) s_\Lambda^{n-1}(r)$  is the surface area element of  $\partial B_R(x)$ . We used the fact  $A(r)$  is non-increasing (Bishop-Gromov inequality) and assumed, without loss of generality,  $R > \sqrt{t}$ .  $\square$

**Lemma 5.5.** *Let  $(M^n, p)$  be a noncompact complete manifold. Then*

$$(5.6) \quad \lim_{R \rightarrow \infty} H_R(x, \cdot, t) = H(x, \cdot, t).$$

*The convergence is uniform, and uniformly in  $L^1$ , on any finite interval  $t \in [0, T]$ .*

*Proof.* Assume  $R > \max\{T, 2d(x, p)\}$ . Put

$$(5.7) \quad M(R) = \sup\{H(x, y, t) | y \in \partial B_R(x), 0 < t \leq T\},$$

by (1.17) and volume comparison we have

$$(5.8) \quad M(R) \leq \sup_{0 < t \leq T} C(n, \Lambda, T) t^{-C_1(n)} e^{-R^2/5t} e^{C_2(n, \Lambda, T)R} \text{Vol}(B_R(p))^{-1},$$

so  $\lim_{R \rightarrow \infty} M(R) \text{Vol}(B_R(p)) = 0$ . By the maximum principle,

$$(5.9) \quad H(x, y, t) - M(R) \leq H_R(x, y, t) \leq H(x, y, t).$$

Combining this with Lemma 5.1, we have

$$(5.10) \quad \|H_R(x, \cdot, t) - H(x, \cdot, t)\|_{L^1} < \epsilon(n, \Lambda, T, R),$$

and  $\lim_{R \rightarrow \infty} \epsilon(n, \Lambda, T, R) = 0$ .  $\square$

Assume  $\lambda_j$  is the  $j$ -th Dirichlet eigenvalue of the Laplacian on  $B_R(p)$ ,  $\phi_j$  is the corresponding eigenfunction,  $\|\phi_j\|_{L^2(B_R(p))} = 1$ .

**Lemma 5.11.** *There exists a constant  $C(n, \Lambda, R)$  such that*

$$(5.12) \quad C(n, \Lambda, R)^{-1} R^{-2} k^{\frac{2}{n}} \leq \lambda_k \leq C(n, \Lambda, R) R^{-2} k^2.$$

*Proof.* Since  $R$  fixed, we have  $\text{Vol}(B_r(x)) \geq C_0(n, \Lambda, R) r^n \text{Vol}(B_R(p))$ , for  $r < 2R$  and  $B_r(x)$  with nonempty intersection with  $B_R(p)$ . Then since  $H_R \leq H$ , we can follow the heat kernel argument as in page 178 of [SY] to get the lower bound of  $\lambda_k$ .

The upper bound follows from an argument of Cheng, see page 105 of [SY].  $\square$

**Lemma 5.13.** *For any  $N > 0$ , there is a function  $\epsilon(N, \Lambda, R, \delta)$  such that for any fixed  $R$ ,  $\lim_{\delta \rightarrow 0} \epsilon(N, \Lambda, R, \delta) = 0$ , And for  $k$  such that  $\lambda_k < N$ ,*

$$(5.14) \quad \int_{A(p, R-\delta, R)} |\phi_k|^2 \leq \epsilon(N, \Lambda, R, \delta).$$

Here  $A(p, R - \delta, R)$  is the annulus  $\{z | R - \delta \leq d(p, z) \leq R\}$ .

*Proof.* By (1.16) and the Bishop-Gromov inequality, when  $t = 1$ ,

$$(5.15) \quad \begin{aligned} \int_{A(p, R-\delta, R)} |\phi_k|^2 &\leq e^{\lambda_k} \int_{A(p, R-\delta, R)} H(x, x, 1) dx \\ &\leq e^N \int_{A(p, R-\delta, R)} \frac{C(n, \Lambda, R)}{\text{Vol}(B_R(p))} dx \leq \epsilon(N, \Lambda, R, \delta). \end{aligned}$$

□

As before, assume  $M_i^n \xrightarrow{d_{GH}} M_\infty$  in the pointed measured Gromov-Hausdorff sense,  $M_i$  is noncompact, satisfies (0.1). Write  $\lambda_{j,i}$  for the  $j$ -th Dirichlet eigenvalue over  $B_R(p_i) \subset M_i^n$ .  $\phi_{j,i}$  is the corresponding eigenvalue:

$$(5.16) \quad \Delta \phi_{j,i} = \lambda_{j,i} \phi_{j,i}; \quad \int_{B_R(p_i)} \phi_{j,i} \phi_{k,i} = \delta_{jk}.$$

**Lemma 5.17.** *For fixed  $j, k > 0$ , assume (for a subsequence of the eigenvalues),  $\lambda_{j,i} \rightarrow \lambda_{j,\infty}$ ,  $\lambda_{k,i} \rightarrow \lambda_{k,\infty}$ . Then there is a subsequence (denoted also by  $\phi_{j,i}, \phi_{k,i}$ ) that converges uniformly on compact subsets of  $B_R$ , and also in  $L^2(B_R)$ , to two locally Lipschitz functions  $\phi_{j,\infty}, \phi_{k,\infty}$ . Moreover,*

$$(5.18) \quad \Delta \phi_{j,\infty} = \lambda_{j,\infty} \phi_{j,\infty}, \quad \Delta \phi_{k,\infty} = \lambda_{k,\infty} \phi_{k,\infty}, \quad \int_{B_R(p)} \phi_{j,\infty} \phi_{k,\infty} = \delta_{jk}.$$

*Proof.* The results is clear in view of Lemma 5.11, Lemma 1.14 and Lemma 3.17. The  $L^2$  convergence and the orthonormal property for the limit functions are implied by locally uniform convergence and Lemma 5.13. □

By Lemma 5.11, we can assume, after passing to a subsequence, that *every* eigenvalue and eigenfunction converge:

$$(5.19) \quad \lim_{i \rightarrow \infty} \lambda_{j,i} = \lambda_{j,\infty}, \quad \lim_{i \rightarrow \infty} \phi_{j,i} = \phi_{j,\infty}.$$

Write

$$(5.20) \quad H_{R,\infty} = \sum_{j=1}^{\infty} e^{-\lambda_{j,\infty} t} \phi_{j,\infty}(x) \phi_{j,\infty}(y).$$

For all fixed  $t, x$ , by Lemma 5.11 and Lemma 1.11, Lemma 1.14,

$$(5.21) \quad H_{R,i}(x, \cdot, t) \rightarrow H_{R,\infty}(x, \cdot, t).$$

The convergence is in  $L^2$ , and is locally uniform. Note we don't know if  $H_{R,\infty}$  (and  $\phi_{j,\infty}, \lambda_{j,\infty}$ ) is well defined. For the moment (before Lemma 5.40), we fix, by a diagonal argument, *one* sequence  $R_k \rightarrow \infty$ , and *one* subsequence  $\{M_{i_v}^n\}$  of  $\{M_i^n\}$  such that for

each  $k$ ,  $H_{R_k,i} \rightarrow H_{R_k,\infty}$ . For simplicity, we just write  $\{M_i^n\}$  for this subsequence of manifolds. So by the results on smooth manifolds, for  $R_j < R_k$ ,

$$(5.22) \quad 0 \leq H_{R_j,\infty}(x, y, t) \leq H_{R_k,\infty}(x, y, t) \leq \frac{C(n, \Lambda)e^{-d^2(x,y)/5t}e^{C\Lambda t}}{\text{Vol}_\infty^{1/2}(B_{\sqrt{t}}(x)) \text{Vol}_\infty^{1/2}(B_{\sqrt{t}}(y))}.$$

Thus we can also assume that the nondecreasing sequence  $H_{R_j,\infty}$  converges pointwise to some function  $H_\infty$ . We will prove that  $H_\infty$  is well defined.

By (5.9) and the locally uniform convergence of  $H_{R,i}$  to  $H_{R,\infty}$  (5.21), the Li-Yau estimate (1.16) is also true for  $H_\infty$ :

$$(5.23) \quad 0 \leq H_\infty(x, y, t) \leq \frac{C(n, \Lambda)e^{-d^2(x,y)/5t}e^{C\Lambda t}}{\text{Vol}_\infty^{1/2}(B_{\sqrt{t}}(x)) \text{Vol}_\infty^{1/2}(B_{\sqrt{t}}(y))}.$$

Note we need to renormalize the measures whenever  $\{M_i^n\}$  is collapsing. Clearly, when  $\text{Ric}_{M_i^n} \geq 0$ , we also have a lower bound of  $H_\infty$  as in (1.18).

**Corollary 5.24.**

$$(5.25) \quad \int_{M_\infty} H_\infty(x, z, s)H_\infty(z, y, t-s)dz = H_\infty(x, y, t).$$

*Proof.* By (5.21), (5.25) is true for  $H_{R,\infty}$ . Write  $H_\infty(x, z, s) = H_{R,\infty}(x, z, s) + \epsilon_R^1(z)$ , similarly  $H_\infty(z, y, t-s) = H_{R,\infty}(z, y, t-s) + \epsilon_R^2(z)$ , here  $H_{R,\infty} = 0$  outside  $B_R(p_\infty)$ ,  $\epsilon_R^1, \epsilon_R^2 \geq 0$  are two functions. In view of Lemmas 5.1, 5.5, (5.21) and (5.23),

$$(5.26) \quad \limsup_{R \rightarrow \infty} (\|\epsilon_R^1(z)\|_{L^1} + \|\epsilon_R^2(z)\|_{L^1}) = 0, \quad \|\epsilon_R^1(z)\|_{L^\infty} + \|\epsilon_R^2(z)\|_{L^\infty} < C(t, s, M_\infty).$$

Now (5.25) is clear. □

**Corollary 5.27.**

$$(5.28) \quad \int_{M_\infty} H_\infty(x, y, t)dy = 1.$$

*Proof.* By (5.21), Lemmas 5.1 and 5.5. □

**Lemma 5.29.** *For any Lipschitz function  $f$  with compact support,*

$$(5.30) \quad \left| \int_{M_\infty} H_{R,\infty}(x, y, t)f(y)dy - f(x) \right| \leq \epsilon(t, \|f\|_{L^\infty}, \mathbf{Lip}f).$$

*Here for any  $F, L > 0$ ,  $\lim_{t \rightarrow 0} \epsilon(t, F, L) = 0$ . The conclusion is also true for  $H_\infty$ .*

*Proof.* By an argument similar to those given in Lemma 5.1 and Lemma 5.5. Note on smooth manifolds, when  $t \rightarrow 0$ , the integral of  $H_R$  is smaller than, but almost equal to 1, and tends to concentrate on smaller and smaller balls centered at  $x$ . In view of (5.21) and the Li-Yau estimate (5.23), we easily get (5.30). □

Let the Sobolev space  $\mathring{H}_{1,2}(B_R(p_\infty))$  be defined as in [Ch3], i.e. the  $H_{1,2}$  closure of the set of Lipschitz functions supported in the interior of  $B_R(p_\infty)$ ,

**Lemma 5.31.** *The space  $\mathring{H}_{1,2}(B_R(p_\infty))$  is contained in  $\Phi$ , the  $L^2$ -linear span of functions  $\phi_{j,\infty}$ . In particular, any Lipschitz function with support in  $B_{R-\delta}$  lies in  $\Phi$ .*

*Proof.* If not, by approximation, we have a Lipschitz function  $f_\infty$  with compact support and an  $\epsilon > 0$  such that

$$(5.32) \quad \sum_{j=1}^{\infty} \left( \int_{B_R(p_\infty)} f_\infty \phi_{j,\infty} \right)^2 < (1 - 3\epsilon) \|f_\infty\|_{L^2}^2.$$

Using Lemma 1.8, we can transplant  $f_\infty$  back to a Lipschitz function,  $f_i$ , on  $M_i^n$ , with compact support which is close to  $f_\infty$  in  $L^\infty$ , such that the energy of  $f_i$  is close to that of  $f_\infty$ . Write

$$(5.33) \quad f_i = \sum_{j=1}^N a_{j,i} \phi_{j,i} + R_{N,i}, \quad R_{N,i} = \sum_{j=N+1}^{\infty} a_{j,i} \phi_{j,i}.$$

Notice,

$$(5.34) \quad \lim_{i \rightarrow \infty} a_{j,i} = \int_{M_\infty} f_\infty \phi_{j,\infty}.$$

So by the min-max principle and Lemma 5.11,  $\lim_{i \rightarrow \infty} \|\nabla f_i\|_{L^2} = \infty$ , we get a contradiction to the construction of  $f_i$ , Lemma 1.8.  $\square$

Remark, it is not clear if we have  $\phi_{j,\infty} \in \mathring{H}_{1,2}$ .

Now for Lipschitz functions  $f_i$  with compact support in  $B_R(p_i) \subset M_i^n$  ( $i = 1, 2, \dots, \infty$ ),  $f_i \rightarrow f_\infty$  uniformly, we have

$$(5.35) \quad f_i = \sum_{j=1}^{\infty} a_{j,i} \phi_{j,i}, \quad a_{j,i} = \int_{M_i^n} f_i \phi_{j,i}.$$

So  $a_{j,i} \rightarrow a_{j,\infty}$ . Clearly,

$$(5.36) \quad \int_{B_R(p_i)} H_{R,i}(x, y, t) f_i(y) dy = \sum_{j=1}^{\infty} e^{-\lambda_j t} a_{j,i} \phi_{j,i}(x).$$

We say  $h(x, t)$  is a *locally strong* solution, if  $h$  continuous, Lipschitz in  $x$ ,  $\frac{\partial h}{\partial t}$  exists, continuous on  $M \times \mathbf{R}^+$ , and when  $t$  fixed,  $-\Delta h = \frac{\partial h}{\partial t}$ , i.e.

$$(5.37) \quad \int_{\Omega} \psi \frac{\partial h}{\partial t} + \int_{\Omega} \langle d_x h, d_x \psi \rangle = 0,$$

for all Lipschitz functions  $\psi$  with compact support.

By Lemma 5.11, Lemma 1.11 and 1.14,

$$(5.38) \quad \lim_{k \rightarrow \infty} \sum_{j=k}^{\infty} |e^{-\lambda_j t} \phi_j(x) d_y \phi_j(y)| = 0.$$

So  $H_{R,i}$  is a locally strong solution of the heat equation. Similarly the function,

$$(5.39) \quad h_i(x, t) = \int_{B_R} H_{R,i}(x, y, t) f_i(y) dy \quad (i = 1, 2, \dots, \infty),$$

is also a locally strong solution. Note for the case  $i = \infty$  we used also Lemma 5.17.

For locally strong solutions on  $M_\infty$ , there is also a weak maximum principle:

**Lemma 5.40.** *Assume  $h$  is a locally strong solution on  $B_{2R} \times [0, T + 1]$ , then if*

$$(5.41) \quad h|_{B_R \times \{0\}} \leq 0, \quad h|_{\partial B_R \times [0, T]} \leq 0.$$

*Then  $h \leq 0$  on  $B_R \times [0, T]$ .*

*Proof.* Define

$$(5.42) \quad m(s) = \sup\{h(x, s) | x \in B_R, \frac{\partial h}{\partial t}(x, s) \leq 0\}.$$

Since  $h, \frac{\partial h}{\partial t}(x, s)$  are continuous functions, it's easy to show that  $m$  is nonincreasing and  $m(0) = 0$  implies  $m(s) \leq 0$  for all  $s > 0$ . Now by the weak maximum principle for Poisson equations (see [GT], [Ch3] or (5.37)), we have, when  $s$  fixed,

$$(5.43) \quad \sup\{h(x, s) | x \in B_R\} = m(s) \leq 0.$$

□

Now we can address the uniqueness of  $H_\infty$ . Recall that  $(M_i, p_i) \rightarrow (M^\infty, p_\infty)$ . Assume for  $R > 0$ , we got two limits  $H_{4R, \infty}^{(1)}, H_{4R, \infty}^{(2)}$  through different subsequences of manifolds.

**Theorem 5.44.** *For  $x, y \in B_R(p_\infty)$ ,  $t < T$ , there is an  $\epsilon(R) > 0$  such that*

$$(5.45) \quad \lim_{R \rightarrow \infty} \epsilon(R) = 0,$$

$$(5.46) \quad H_{4R, \infty}^{(1)}(x, y, t) < H_{4R, \infty}^{(2)}(x, y, t) + \epsilon(R).$$

*Proof.* We can assume  $R > T^2 > t^2$  and  $R > 4$ . Assume (5.46) is not true, then there is a point  $a \in B_R(p_\infty)$  and  $0 < r < 1$  such that

$$(5.47) \quad H_{4R, \infty}^{(1)}(x, y, t) \geq H_{4R, \infty}^{(2)}(x, y, t) + \epsilon(R),$$

for  $y \in B_{2r}(a)$ . We then construct a test function  $f \geq 0$  such that,  $f$  Lipschitz, supported in  $B_r(a)$ ,

$$(5.48) \quad 2 \int_{B_r(a)} f \geq \text{Vol}(B_r(a)) \sup_{B_r(a)} f.$$

Clearly, for  $R < \infty$ , the functions,

$$(5.49) \quad F_k(z, s) = \int_{B_r(a)} H_{4R, \infty}^{(k)}(z, y, s) f(y) dy, \quad (k = 1, 2),$$

are locally strong solutions of the heat equation, and (by the construction of  $f$ ),

$$(5.50) \quad F_1(x, t) \geq F_2(x, t) + \epsilon \int_{B_r(a)} f \geq F_2(x, t) + \frac{\epsilon(R)}{2} \text{Vol}(B_r(a)) \sup_{B_r(a)} f.$$

For a point  $z$  near  $\partial B_{2R}(p_\infty)$ , say  $d(z, p) = 2R$ ,  $d(a, z) \geq R$ ,

$$(5.51) \quad F_k(z, s) < \text{Vol}(B_r(a)) \frac{C(n)}{\text{Vol}(B_{\sqrt{s}}(z))} e^{-R^2/5s} e^{CR} \sup_{B_r(a)} f, \quad (k = 1, 2).$$

By a standard argument of the Bishop-Gromov inequality,

$$(5.52) \quad F_k(z, s) < C_1(n) \frac{\text{Vol}(B_r(a))}{\text{Vol}(B_1(p))} s^{-\frac{n}{2}} e^{-R^2/5s} e^{C'R} \sup_{B_r(a)} f, \quad (k = 1, 2).$$

Next we consider the case that  $s$  is small. Since  $f$  is fixed, by (5.30),  $F_k \rightarrow f$  uniformly on  $B_{2R}(p)$  when  $s \rightarrow 0$ .

In view of the weak maximum principle on  $B_{2R}(p_\infty) \times [0, T]$  (Lemma 5.40), clearly we should choose  $\epsilon(R)$  such that that for  $0 \leq s \leq T$ ,

$$(5.53) \quad \frac{C_1(n)}{\text{Vol}(B_1(p))} s^{-\frac{n}{2}} e^{-R^2/5s} e^{C'R} < \frac{\epsilon(R)}{4},$$

by the maximum principle we got a contradiction to (5.50).  $\square$

**Theorem 5.54.**  *$H_\infty$  is well defined. For fixed  $t > 0$ ,  $x_i \rightarrow x_\infty$ , we have  $H_i(x_i, \cdot, t) \rightarrow H_\infty(x_\infty, \cdot, t)$  in  $L^1$ . When  $H_\infty$  is continuous, this convergence is also uniform.*

*Proof.* By the previous theorem and the construction of  $H_\infty$  (compare (5.9)), we see  $H_\infty$  is independent of the choice of subsequences, so well defined.

We already know, by (5.9), (5.21), (5.23), that locally  $H_i \rightarrow H_\infty$  in  $L^1$ . The proof of global  $L^1$  convergence is similar with Lemma 5.1, Lemma 5.5, using (1.18), (5.23).

Recall (see [SY] Chapter 4), there is a Harnack inequality

$$(5.55) \quad H_i(x, y_1, t_1) \leq H_i(x, y_2, t_2) \left(\frac{t_2}{t_1}\right)^n \exp\left(\frac{d^2(y_1, y_2)}{4(t_2 - t_1)} + C(n, \Lambda)(t_2 - t_1)\right),$$

for  $0 < t_1 < t_2$ . If  $H_\infty$  is continuous, then locally  $H_\infty$  is uniformly continuous (especially, with respect to  $t$ ), clearly by (5.55) the convergence  $H_i \rightarrow H_\infty$  must be uniform, compare with (5.23).  $\square$

We now want to interpret the meaning of  $H_\infty$ . Recall from [ChCo4] and [Ch3],  $\Delta$  is a *positive self-adjoint operator*. So  $-\Delta$  generates a semigroup  $e^{-t\Delta}$ .

Assume  $f_i$  is supported in  $B_K(p_i) \subset B_R(p_i)$ . Use the notation in (5.35), define

$$(5.56) \quad W_{R,i}(t)f_i(x) = \sum_{j=1}^{\infty} a_{j,i} \cos(\sqrt{\lambda_{j,i}}t)\phi_{j,i}.$$

By the finite speed of propagation (see [Ta]), when  $t$  is fixed and  $R > K + t$ ,  $W_{R,i}(t)f$  is independent of  $R$ . We write  $W_i(t)f$  for  $W_{R,i}(t)f$  with  $R$  big. For  $i < \infty$ ,

$$(5.57) \quad e^{-t\Delta}f_i(x) = \int_{M_i^n} H_i(x, y, t)f_i(y)dy = \int_{-\infty}^{\infty} e^{-s^2/4t}W_i(s)f_i(x)ds,$$

see [CGT], [Ta]. Define

$$(5.58) \quad W_{R,\infty}(t)f_\infty(x) = \sum_{j=1}^{\infty} a_{j,\infty} \cos(\sqrt{\lambda_{j,\infty}}t)\phi_{j,\infty}.$$

We notice that  $W_{R,i}$  ( $i = 1, 2, \dots, \infty$ ) does not increase  $L^2$  norm, and we should use Lemma 1.8 and approximation to construct  $C^2$  functions  $f_i$  on  $M_i^n$  that converges to  $f_\infty$ . Clearly,  $W_{R,i}f_i \rightarrow W_{R,\infty}f_\infty$  in  $L^2$ . We remark that generally, we don't know if  $W_{R,\infty}$  is well defined.

**Theorem 5.59.** *If the limit  $M_\infty$  is a smooth manifold, and the limit measure is the canonical measure on  $M_\infty$ , then  $H_\infty$  is the heat kernel on  $M_\infty$ .*

*Proof.* In the noncollapsing case, by Colding's theorem [Co], the limit measure is the canonical measure on  $M_\infty$ ; when  $M_\infty = \mathbf{R}^k$  for some  $k$ , the limit measure is also a multiple of the standard Lebesgue measure on  $\mathbf{R}^k$ , see [ChCo2]. In these cases, the Laplacian we defined on  $M_\infty$  is the same one from the original smooth structure of  $M_\infty$ .

Pick any  $C_0^\infty$  function  $f$  supported in  $B_R$ , So

$$(5.60) \quad \int_{M_\infty} (\Delta^k) f \phi_{j,\infty} = (\lambda_{j,\infty})^k \int_{M_\infty} f \phi_{j,\infty} = (\lambda_{j,\infty})^k a_{j,\infty}.$$

Since  $(\Delta^k) f \in C_0^\infty$ , we have for all  $k$ ,  $\lim_{j \rightarrow \infty} (\lambda_{j,\infty})^k a_{j,\infty} = 0$ . By Lemma 5.11, we have for all  $k$ ,  $\lim_{j \rightarrow \infty} j^k a_{j,\infty} = 0$ . So  $W_{R,\infty}(t)f$  is a classical solution of the wave equation, when  $R$  is big enough,  $W_{R,\infty}(t)f = W_\infty(t)f$  is independent of  $R$ . Since  $M_\infty$  is a smooth manifold,

$$(5.61) \quad e^{-t\Delta} f(x) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-s^2/4t} W_\infty(s) f(x) ds.$$

In view of (5.57), combined with the fact  $W_\infty$  does not increase  $L^2$  norm and  $H_i(x, y, t)$  converges uniformly to  $H_\infty(x, y, t)$ , we have

$$(5.62) \quad e^{-t\Delta} f(x) = \int_{M_\infty} H_\infty(x, y, t) f(y) dy.$$

That concludes the proof.  $\square$

## 6. LAPLACIAN ON METRIC CONES

In this section, we assume  $M_i^n \xrightarrow{d_{GH}} C(X)$  where  $C(X)$  is a metric cone;  $\text{Ric}_{M_i^n} \geq 0$ ,  $M_i^n$  is complete noncompact and satisfies (0.4) uniformly,  $n \geq 3$ .

Write  $p_\infty$  for the pole of  $C(X)$ , define  $r(x) = d(x, p_\infty)$ .

**Theorem 6.1.** *If  $M_\infty = C(X)$ , then  $H_\infty$  is the integral kernel of the semigroup  $e^{-t\Delta}$ .*

*Proof.* In view of (5.23), (5.25) and the Young's inequality, one can define a semigroup  $E(t)$  on  $L^2(M_\infty)$  by

$$(6.2) \quad E(t)f(x) = \int_{M_\infty} H_\infty(x, y, t) f(y) dy.$$

We want to compare  $E(t)$  with  $e^{-t\Delta}$ . First, by Theorem 3.21, (1.18) and (5.23), one easily get

$$(6.3) \quad G_\infty(x, y) = \int_0^\infty H_\infty(x, y, t) dt.$$

Pick any  $L^2$  function  $f$  with compact support. Write

$$(6.4) \quad F(x) = \int_{M_\infty} G_\infty(x, y) f(y) dy.$$

We compute

$$\begin{aligned}
(6.5) \quad \frac{E(t)F - F}{t} &= \int_{M_\infty} \left( \frac{H_\infty(x, y, t)}{t} \right) \int_{M_\infty} \int_0^\infty H_\infty(y, z, s) f(z) ds dz dy \\
&\quad - \frac{1}{t} \int_{M_\infty} \int_0^\infty H_\infty(x, z, s) f(z) ds dz \\
&= -\frac{1}{t} \int_0^t \int_{M_\infty} H_\infty(x, z, s) f(z) dz ds.
\end{aligned}$$

So by (0.4), (5.23), (5.30) and the Young's inequality we have

$$(6.6) \quad \lim_{t \rightarrow 0} \frac{E(t)F - F}{t} \rightarrow -f.$$

in  $L^2$  and  $L^1$ .

Now we use the assumption that  $M_\infty = C(X)$  is a noncollapsed cone. Recall the results in Section 4, we can construction a function  $\phi = \phi(r)$  such that  $\phi$  is a smooth function of  $r$ , where  $r(x) = d(p_\infty, x)$  is the distance from the pole, and

$$(6.7) \quad \phi(r) = 1 \text{ if } r < R, \quad \phi(r) = 0 \text{ if } r \geq R + 2, \quad \nabla \phi \leq C_0 \sqrt{\phi}.$$

So on  $M_\infty = C(X)$  we have  $\Delta \phi = -\phi'' - (n-1)\phi'/r$ . This function can serve as a cut off function.

We prove, if  $F, f = \Delta F \in L^2$  have compact support, then

$$(6.8) \quad F = \int_{C(X)} G_\infty(x, y) f(y) dy.$$

In fact, assume  $\{f_k\}$  is a sequence of Lipschitz functions,  $f_k \rightarrow f$  in  $L^2$ , and all  $f_k$  together with  $f, F$  are supported in the ball  $B_K(p_\infty)$ . So the function

$$(6.9) \quad F_k = \int_{C(X)} G_\infty(x, y) f_k(y) dy,$$

satisfies  $\Delta F_k = f_k$  by the discussion in Section 3. Consider the equation  $\Delta(F_k - F) = f_k - f$ , i.e.

$$(6.10) \quad \int_{C(X)} \langle dF_k - dF, du \rangle - \int_{C(X)} (f_k - f)u = 0,$$

for any  $u \in \mathring{H}_{1,2}$ . We set  $u = \phi(F_k - F)$ , so  $du = d\phi(F_k - F) + \phi(dF_k - dF)$ . By the Schwartz inequality,

$$(6.11) \quad \begin{aligned} &\|\sqrt{\phi}d(F_k - F)\|_{L^2}^2 - C_0 \|(F_k - F)|_{A(R, R+2)}\|_{L^2} \|\sqrt{\phi}d(F_k - F)|_{A(R, R+2)}\|_{L^2} \\ &\quad - \|(f_k - f)|_{B_K}\|_{L^2} \|(F_k - F)|_{B_K}\|_{L^2} \leq 0, \end{aligned}$$

here  $A(R, R+2)$  is the annulus  $\{x | R \leq r(x) \leq R+2\}$ . Note we have a definite bound for  $\|F_k|_{B_K}\|_{L^2}$  by (1.19) and the Young's inequality. Note also by (1.19) we get, for  $R > K$ ,

$$(6.12) \quad \begin{aligned} \|(F_k - F)|_{A(R, R+2)}\|_{L^2} &= \|F_k|_{A(R, R+2)}\|_{L^2} < C(n, \|f\|_{L^1}) (R^{4-2n} R^{n-1})^{1/2} \\ &= C(n, \|f\|_{L^1}) R^{(3-n)/2} < C(n, \|f\|_{L^1}), \end{aligned}$$

since  $n \geq 3$ . So first, we get that  $\|d(F_k - F)\|_{L^2} < \infty$  by letting  $R \rightarrow \infty$ . Then by letting  $k \rightarrow \infty$ , we have  $\|d(F_k - F)\|_{L^2} \rightarrow 0$ , since we can choose  $R$  in (6.11) such that  $\|\sqrt{\phi}d(F_k - F)|_{A(R, R+2)}\|_{L^2}$  small.

Now by the (2, 2)-Poincare inequality, (0.4), (1.19) and Young's inequality,  $F_k \rightarrow F$  in  $L^2$  on compact sets. Also notice, on any compact sets, the right hand side of (6.9) converges to the right hand side of (6.8) in  $L^2$ , by the Young inequality (however, in view of (1.19), these convergences might not be globally  $L^2$ ). That's enough to imply (6.8).

Next we compute, for  $f = \Delta F$ ,  $F, f \in L^2$ ,

$$(6.13) \quad \begin{aligned} \|\Delta(\phi F) - f\|_{L^2} &\leq \|F\Delta\phi\|_{L^2} + \|(\phi - 1)f\|_{L^2} + 2\|d\phi, dF\|_{L^2} \\ &\leq C(n)\|F|_{A(R, R+2)}\|_{L^2} + \|f|_{A(R, R+2)}\|_{L^2} + C_0\|dF|_{A(R, R+2)}\|_{L^2}. \end{aligned}$$

Similar to (6.11), one shows  $\|dF\|_{L^2} < \infty$ . So if  $R \rightarrow \infty$ , we have  $\phi F \rightarrow F$  and  $\Delta(\phi F) \rightarrow f = \Delta F$  in  $L^2$ . Moreover, by (6.8),

$$(6.14) \quad \phi F(x) = \int_{C(X)} G_\infty(x, y)\Delta(\phi F)(y)dy.$$

So the computation (6.5) is valid for the functions  $\phi F$  and  $\Delta(\phi F)$ :

$$(6.15) \quad \lim_{t \rightarrow 0} \frac{E(t)\phi F - \phi F}{t} = -\Delta(\phi F)$$

in  $L^2$ . We already know  $E(t)$  in (6.2) is a semigroup, its infinitesimal generator is a *closed* operator (see [Ta]). So by the above computations, this infinitesimal generator must be the self-adjoint operator  $-\Delta$  on  $C(X)$ .  $\square$

By the discussion in the beginning of Section 4, we have an eigenfunction expansion of Laplacian on the unit cross section  $X$ . We denote by  $\phi_j$  ( $j = 0, 1, 2, \dots$ ) the renormalized eigenfunctions with eigenvalues  $\mu_j > 0$ , note  $\phi_0 = \text{Vol}(X)^{-1/2}$ .  $\mu_j \rightarrow \infty$  when  $j \rightarrow \infty$ .

Put  $d = \text{Diam } X$ . Using an argument of Gromov (see [Gr], and Theorem 4.8 of [Ch3]), we have a more precise estimate of  $\mu_j$ :

$$(6.16) \quad \mu_j > C(\tau, \kappa)^{-1}d^{-2}j^{\frac{2}{\kappa}}.$$

On the other hand, on each ball  $B_r(x_k)$  of radius  $r = d/2(j+2)$  on  $X$ , we define a Lipschitz function  $\psi_k$  supported in  $B_r(x_k)$  using MacShane's lemma ([Ch3], [ChCo3]):

$$(6.17) \quad \psi_k(x_k) = r, \quad \psi_k(\partial B_r(x_k)) = 0, \quad \mathbf{Lip}\psi_k = 1,$$

so we can follow the argument of Cheng (see p.105 of [SY]), and get

$$(6.18) \quad \mu_j \leq C(\kappa)j^2d^{-2}.$$

Now we can use Moser iteration,  $|\phi_j|$  is bounded by a definite power of  $j$ :

$$(6.19) \quad |\phi_j| \leq C(d, \kappa, \tau)j^{N(\tau, \kappa)}.$$

Moreover,  $\phi_j$  is Hölder continuous, see [GT], [Lin].

Write  $\nu_j = \sqrt{\mu_j + \alpha^2}$ , here  $m = n - 1$ ,  $\alpha = (1 - m)/2$ . We write  $x, y$  in polar coordinates,  $x = (r_1, x_1), y = (r_2, x_2)$ .

**Theorem 6.20.**

$$(6.21) \quad H_\infty = (r_1 r_2)^\alpha \sum_{j=0}^{\infty} \left(\frac{1}{2t}\right) e^{-(r_1^2 + r_2^2)/4t} I_{\nu_j} \left(\frac{r_1 r_2}{2t}\right) \phi_j(x_1) \otimes \phi_j(x_2).$$

Here  $I_{\nu_j}$  are the modified Bessel functions:

$$(6.22) \quad I_\nu(z) = \left(\frac{z}{2}\right)^\nu \sum_{k=0}^{\infty} \frac{1}{k! \Gamma(\nu + k + 1)} \left(\frac{z}{2}\right)^{2k}.$$

In our case  $\Delta$  is a self-adjoint operator on the whole cone  $C(X)$ , namely, including the pole  $p_\infty$ . By Corollary 4.25, the separation of variable formula (4.16) works for  $u = f(r)g(x)$  on the whole  $C(X)$  if  $u$  and  $\Delta u$  are bounded on  $C(X) \setminus \{p_\infty\}$ . So the heat kernel on  $M_\infty$  has the expression as on the right hand side of (6.21); the proof goes exactly like the classical case, see [Ch1], [Ch2] page 592, [ChTa1] and [Ta] chapter 8, we omit the details. By Theorem 6.1, we have (6.21).

By Stirling's formula, (6.16) and (6.18), we see the series (6.21) converges uniformly, when  $t$  is bounded away from 0 and  $r_1, r_2$  stay bounded. In particular,  $H_\infty$  is continuous, so by Theorem 5.54 we have  $H_i \rightarrow H_\infty$  uniformly.

If one of the two points  $x$  and  $y$ , say,  $y$ , is the pole  $p_\infty$ , then there is only one term in (6.21). Note  $\nu_0 = -\alpha = (m-1)/2$ ,  $m = n-1$ ,

$$(6.23) \quad H_\infty(p_\infty, x, t) = \left(\frac{1}{4\pi t}\right)^{\frac{n}{2}} e^{-(r^2)/4t} \frac{2\pi^{n/2}}{\Gamma(n/2)} (\text{Vol}(X))^{-1}.$$

As a corollary, we get a new proof of Li's asymptotic formula for heat kernels [Li1]:

**Corollary 6.24** (Li). *Assume  $M^n$  is a complete noncompact manifold satisfying (0.4),  $\text{Ric}_{M^n} \geq 0$ . Then*

$$(6.25) \quad \lim_{t \rightarrow \infty} \text{Vol}(B_{\sqrt{t}}(p)) H(p, y, t) = (4\pi)^{-n/2} \omega_n.$$

$\omega_n$  is the volume of the unit ball in  $\mathbf{R}^n$ .

*Proof.* Notice,

$$(6.26) \quad \lim_{t \rightarrow \infty} \text{Vol}(B_{\sqrt{t}}(p)) t^{-n/2} = v_0 = n^{-1} \text{Vol}(X).$$

So we need to show,

$$(6.27) \quad \lim_{t \rightarrow \infty} t^{n/2} \text{Vol}(X) H(p, y, t) = (4\pi)^{-n/2} n \omega_n.$$

Assume  $t_i \rightarrow \infty$ ,  $M_i^n = (M^n, p, t_i^{-1} dx^2) \xrightarrow{d_{GH}} C(X)$  for some metric cone  $C(X)$ ; see [ChCo1]. The heat kernel  $H_i(p, x, t)$  on  $M_i^n$  is

$$(6.28) \quad H_i(p, y, 1) = t_i^{n/2} H(p, y, t).$$

Here we identify  $p, x \in M_i^n$  with  $p, x \in M$ , however,  $d_{M_i^n}(p, x) = t_i^{-1/2} d_M(p, x)$ ,  $d_{M_i^n}$  is the distance on  $M_i^n$ . In particular,  $d_{M_i^n}(p, x) \rightarrow 0$  as  $i \rightarrow \infty$ . Since  $M_i^n \xrightarrow{d_{GH}} C(X)$ ,

by Theorem 5.54 and (6.23) we have

$$(6.29) \quad \begin{aligned} \lim_{t \rightarrow \infty} t^{n/2} \text{Vol}(X)H(p, y, t) &= \text{Vol}(X) \lim_{i \rightarrow \infty} H_i(p, x, 1) \\ &= \text{Vol}(X)H_\infty(p_\infty, p_\infty, 1) = (4\pi)^{-n/2} \frac{2\pi^{n/2}}{\Gamma(n/2)}. \end{aligned}$$

We just need to recall  $n\omega_n = 2\pi^{n/2}(\Gamma(n/2))^{-1}$  (see [Ta] Chapter 3).

Finally in view of the almost rigidity theorem [ChCo1], we see the above results holds for *all* sequences  $t_i \rightarrow \infty$ . This suffices to complete the proof.  $\square$

Similarly, we get the asymptotic formula for heat kernels in [LiTW]:

**Corollary 6.30** (Li-Tam-Wang). *Assume  $M^n$  is a complete noncompact manifold satisfying (0.4),  $\text{Ric}_{M^n} \geq 0$ . Then for  $p \in M^n$ , and any  $R, T > 0$ ,*

$$(6.31) \quad \lim_{d(p,x) \rightarrow \infty} \text{Vol}(B_{R^{-1}d(p,x)}(p))H(p, x, Td(p,x)^2R^{-2}) = \frac{\omega_n}{(4\pi T)^{n/2}e^{R^2/4T}}.$$

*Proof.* We use the same argument as in Corollary 6.24. For  $x_i$  with  $d(p, x_i) \rightarrow \infty$ , we study the heat kernels on the sequence  $M_i^n = (M^n, p, R^2d(p, x_i)^{-2}dx^2)$ .  $\square$

We can similarly get a *local* asymptotic formula for  $H_\infty$ .

## 7. EIGENVALUES ON COMPACT LIMIT SPACES

We assume  $M_i^n \xrightarrow{dGH} M_\infty$ , with  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$ ,  $M_\infty$  compact. A point  $x \in M_\infty$  is said to be regular,  $x \in \mathcal{R}_k$ , if all tangent cones at  $x$  equal to  $\mathbf{R}^k$ ; see [ChCo2].

**Lemma 7.1.** *If  $x \in \mathcal{R}_n \subset M_\infty$ , then*

$$(7.2) \quad \lim_{t \rightarrow 0} H_\infty(x, x, t)t^{\frac{n}{2}} = (4\pi)^{-\frac{n}{2}}.$$

*Proof.* Use a similar argument as the one in Corollary 6.24.  $\square$

**Theorem 7.3.** *Assume  $M_i^n \xrightarrow{dGH} M_\infty$ ,  $\text{Ric}_{M_i^n} \geq -(n-1)\Lambda$ , and for some  $v_0 > 0$ ,  $\text{Vol}(M_i^n) \geq v_0$ . Then*

$$(7.4) \quad \lim_{j \rightarrow \infty} j^{-\frac{2}{n}} \lambda_{j,\infty} = 4\pi\Gamma\left(\frac{n}{2} + 1\right)^{\frac{2}{n}} \mu_\infty(M_\infty)^{-\frac{2}{n}}.$$

*Proof.* In this case we don't need to renormalize the volume on  $M_i^n$  (see [ChCo2]). Note for some  $D$  we have  $\text{Diam } M_i^n \leq D$ ,  $i = 1, 2, \dots, \infty$ , by the Bishop-Gromov inequality and (1.16), we get

$$(7.5) \quad t^{\frac{n}{2}} H_\infty(x, x, t) \leq C(n, \Lambda, D, v_0).$$

Moreover, almost every point of  $M_\infty$  is in  $\mathcal{R}_n$ . Now by Corollary 7.1, for  $x \in \mathcal{R}_n$ ,  $t^{\frac{n}{2}} H_\infty(x, x, t) \rightarrow (4\pi)^{-n/2}$  when  $t \rightarrow 0$ . By the dominated convergence theorem,

$$(7.6) \quad \lim_{t \rightarrow 0} t^{\frac{n}{2}} \int_{M_\infty} H_\infty(x, x, t) dx = (4\pi)^{-\frac{n}{2}} \mu_\infty(M_\infty).$$

Finally, by applying the Karamata Tauberian theorem (see [Ta] Chapter 8), we have

$$(7.7) \quad \lim_{\lambda \rightarrow \infty} \lambda^{-\frac{n}{2}} N(\lambda) = \mu_\infty(M_\infty) \Gamma\left(\frac{n}{2} + 1\right)^{-1} (4\pi)^{-\frac{n}{2}},$$

where  $N(\lambda)$  is the number of eigenvalues smaller than  $\lambda$ . Clearly this implies the Weyl asymptotic formula (7.4).  $\square$

When the limit space  $M_\infty$  is collapsed, at present our results are less satisfactory. Recall the notion of *Minkowski dimensions*; see [Ma]. Assume  $Z$  is a metric space. For  $d > 0$ , let  $N(Z, \epsilon) \in \mathbf{Z}$  be the minimal integer such that  $Z$  can be covered by  $N(Z, \epsilon)$  many balls of radius  $\epsilon$ . Put

$$(7.8) \quad v_d^-(Z) = \liminf_{\epsilon \rightarrow 0} \epsilon^d N(Z, \epsilon),$$

$$(7.9) \quad v_d^+(Z) = \limsup_{\epsilon \rightarrow 0} \epsilon^d N(Z, \epsilon).$$

Here  $v_d^\pm(M_\infty)$  can be  $\infty$ . The upper (lower) Minkowski dimension is defined by

$$(7.10) \quad \overline{\dim}_{Mink}(Z) \ (\underline{\dim}_{Mink}(Z)) = \inf\{d | v_d^+(Z) = 0 \ (v_d^-(Z) = 0)\}.$$

**Lemma 7.11.** *There exist  $E_1(n), E_2(n) > 0$  such that for any  $d > 0$ ,*

$$(7.12) \quad \limsup_{t \rightarrow 0} t^{\frac{d}{2}} \int_{M_\infty} H_\infty(x, x, t) dx \leq E_2 v_d^+(M_\infty),$$

and if, in addition,  $\text{Ric}_{M_i^n} \geq 0$ , then

$$(7.13) \quad E_1 v_d^-(M_\infty) \leq \liminf_{t \rightarrow 0} t^{\frac{d}{2}} \int_{M_\infty} H_\infty(x, x, t) dx.$$

*Proof.* Let  $\cup_{1 \leq j \leq N(M_\infty, \sqrt{t})} B_{\sqrt{t}}(x_j)$  be a covering of  $M_\infty$  by a minimal set of balls of radius  $\sqrt{t}$ . We add up the integrals of  $H_\infty$  on these ball an use Corollary 2.7 to get the estimates (7.12), (7.13).  $\square$

**Lemma 7.14.** *If  $v_d^+(M_\infty) < c < \infty$ , then there exist  $C$  such that*

$$(7.15) \quad \lambda_{j, \infty} > C j^{\frac{2}{d}}.$$

*Proof.* We can follow an argument of Gromov (see [Gr] or Theorem 4.8 in [Ch3]). Here we use the assumption  $v_d^+(M_\infty) < c < \infty$  to estimate the number of balls that is needed to cover  $M^\infty$ .  $\square$

**Lemma 7.16.** *If  $v_d^-(M_\infty) > c > 0$ , then there exist  $C$  depending on  $n, c$ , such that*

$$(7.17) \quad \lambda_{j, \infty} \leq C j^{\frac{2}{d}}.$$

*If  $k$  is the maximal integer such that  $\mathcal{R}_k \subset M_\infty$  is not empty, then*

$$(7.18) \quad \lambda_{j, \infty} < C(M_\infty)(j)^{\frac{2}{k}}.$$

*Proof.* For  $r > 0$ ,  $M_i^n$  contains  $j = C(n, c)r^{-d}$  many disjoint balls of radius  $r$  for  $i$  big enough. The result follows by a well known argument of Cheng [Cheng]; see page 105 of [SY].

If  $k$  is the maximal integer such that  $\mathcal{R}_k \subset M_\infty$  is not empty, then the  $k$ -Hausdorff measure of  $M_\infty$  is positive (see [ChCo3] or [Ch3]). So  $v_k^-(X) > 0$ . By (7.17) we get (7.18).  $\square$

If one can also prove for any  $d > k$ ,

$$(7.19) \quad \lim_{t \rightarrow 0} t^{\frac{d}{2}} \sum_{j=0}^{\infty} e^{-\lambda_{j,\infty} t} = \lim_{t \rightarrow 0} t^{\frac{d}{2}} \int_{M_\infty} H_\infty(x, x, t) dx = 0,$$

then by Lemma 7.11,  $d_M(M_\infty)$ , the Minkowski dimension of  $M_\infty$  is no more than  $k$ . Combine with the results in [ChCo3] and [Ch3],  $d_M(M_\infty) = k$ . However, at present we don't know how to get (7.19). One related question is,

*Question:* Is there an  $\epsilon(n) > 0$ , such that for any  $M^n$  with  $\text{Ric}_{M^n} \geq 0$ , any eigenfunction  $\phi$  of  $\Delta$  and any set  $E$  with  $\text{Vol}(E) < \epsilon \text{Vol}(M)$ , we have

$$(7.20) \quad \int_{M^n - E} \phi^2 > \frac{1}{2} \int_{M^n} \phi^2 ?$$

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