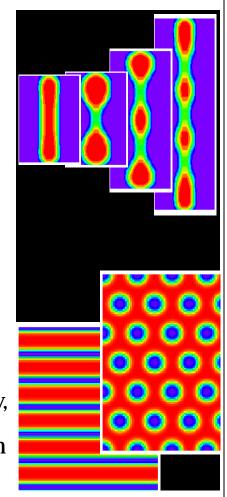


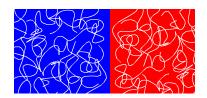
# Self-Consistent Field Methods in Polymer Physics

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- Stuff we want to do:
- Strengthen mixtures of plastics

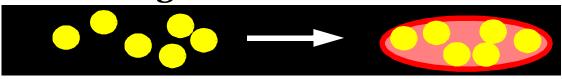


- Incompatible plastics
- Combine properties (strength, flexibility)
- Lubricate/protect surfaces

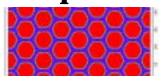


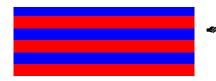
- Prevent contact
- Avoid damage

Encapsulate drugs



Create patterns





Symmetry, scale

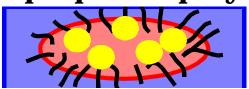
- Stuff that can do it.
- **Stitching polymers: reinforce mixtures**



- Half blue/half red reinforces interface.
- **End-grafted polymers: lubrication**



- Trapped coating "Osmotic" barrier
- **Amphiphillic polymers: housing for droplets**



- Polymer forms vesicles
- Release contents, pH e.g.
- **Block copolymers: templates for ordering**

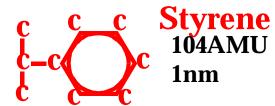




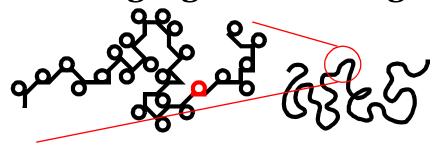
http://www.princeton.edu/~polymer/

#### Polymers

□ Are made of monomers...



... strung together into huge chains...



# **Polystyrene**

1000 monomers: 104,000AMU  $1000 nm = 1 \mu m$ 

 $\Box$  ... which mostly ignore h...

$$\Delta x \Delta p \approx (1 nm)(10^5 \text{AMU} v) \approx 10^8 \frac{v}{m/s} h$$

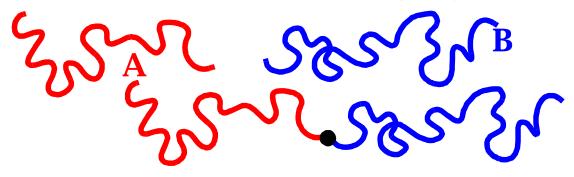
... and are all tangled up.



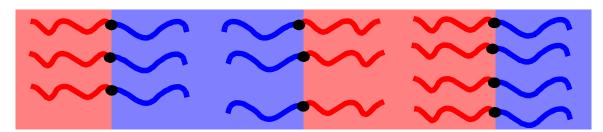
Entangling, knots

## Block copolymers

**□** Two kinds of monomers strung together.

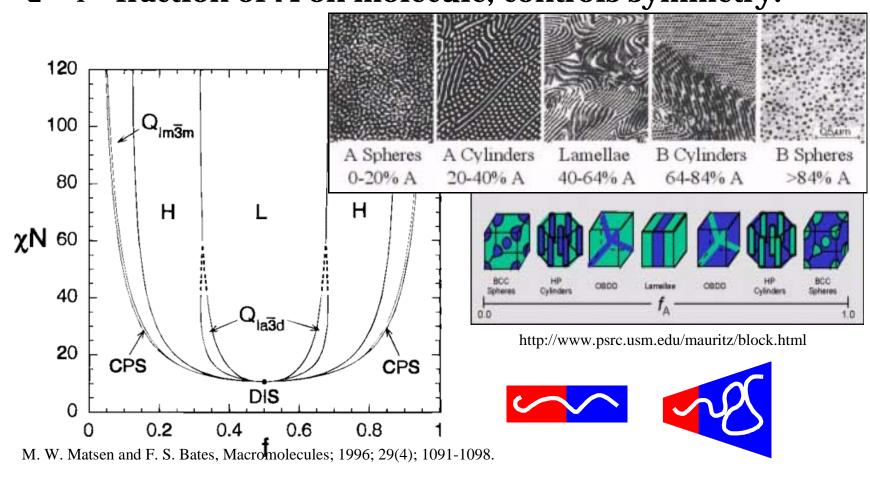


- A-block and B-block: "diblock"
- Unless you break bonds, micro-scale texture happens.

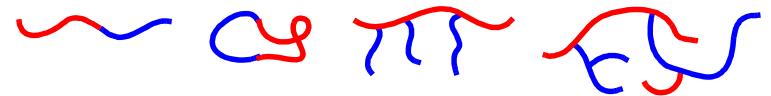


# Asymmetric diblocks

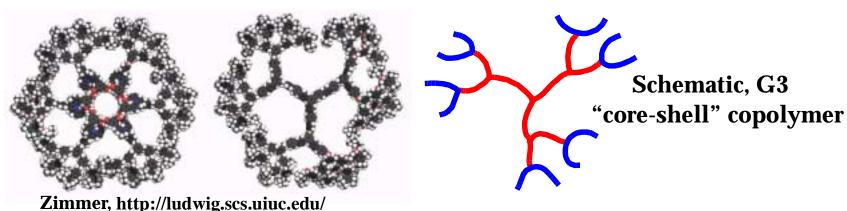
 $\Box$  *f*= fraction of A on molecule, controls symmetry:



- Architecture controls properties
- Diblocks, composition fraction is only control
- Chain topology is also something to consider:

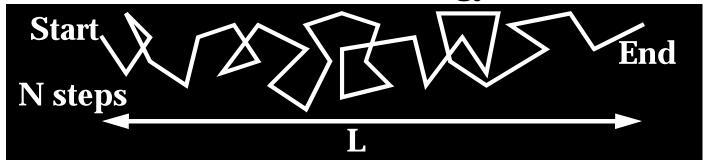


"Starburst" dendrimers



- Outline
- Self-consistent field treatment of complex polymer mixtures
- Applied to dendrimers
- Applied to charged polymers
- Conclusions

- **Self-Consistent Field**
- Random walk, N steps covering an end-end distance L. How much free energy to bias it?



Apply a force F to each link.

$$force \propto L$$

$$force \propto 1/N$$

Double F, go twice as far More steps, less force

- Work = force \*distance:  $\left(\frac{L}{N}\right)L = \left(\frac{L}{N}\right)^2 N$ In all:  $S = \int dn \left(\frac{dl}{dn}\right)^2$ , non-uniform stretching.

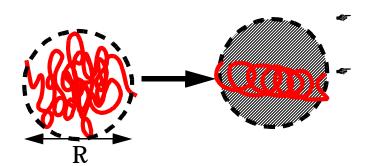
- Random walk in a field
- Field biases walk, statistics:



- □ Walks are biased toward lower values of monomer chemical potential: U(r)
- □ Choose U(r) to match some property of the surrounding material
  - E.g.  $U(r) \approx \varphi(r)$ , monomer volume fraction, polymer solution.
  - $\phi(r) \approx \text{ probability that one of the surrounding biased walks has a monomer at } r$

#### One polymer by itself, excluded volume

#### Random walk in a potential:



Chain crossings result in an outward pressure

Chain entropy pulls surface in: effective surface tension:

$$\gamma \approx \frac{(R/N)}{R}$$

Balance:

$$\frac{\gamma}{R} = \Delta P$$

#### $\Box$ Relate $\triangle P$ to polymer properties:

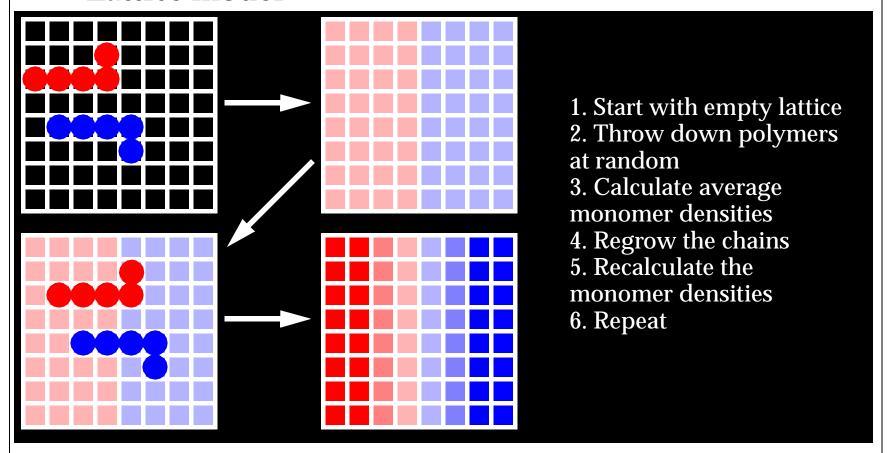
2nd virial approach gives SAW scaling of Flory:

$$\Delta P = c^2 = \frac{N^2}{R^6} \longrightarrow \frac{1}{R} \frac{1}{N} = \frac{N^2}{R^6} \longrightarrow R = N^{3/5}$$

#### Numerical self-consistent lattice calculations

Fleer, Cohen, Scheutjens, Cosgrove, Vincent, Polymers at Interfaces Chapman and Hall, London 1993

Lattice model



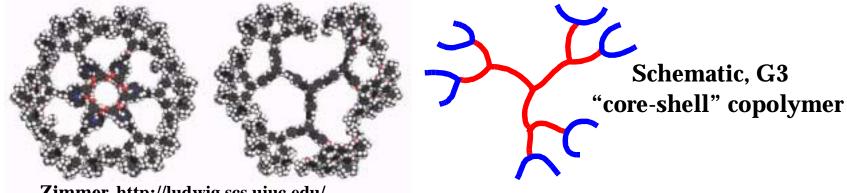
Dendrimers, charged blend, lattice electrostatics.

Outline

- Self-consistent field treatment of complex polymer mixtures
- Applied to Dendrimers
- Applied to charged polymers
- **Conclusions**

#### **Dendrimers**

Proliferation of tips on a single molecule:



- Zimmer, http://ludwig.scs.uiuc.edu/
- Complex self-organization for a single molecule
- Applications depend on
  - Monomer density
  - Location of tips

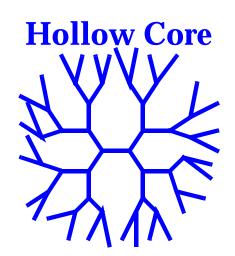
#### Hollow or Filled Core?

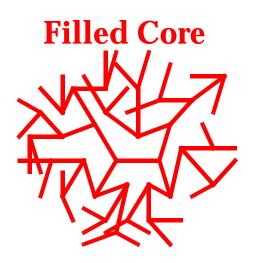
#### Hervet and deGennes

- Long, flexible spacers
- Tips segregate spontaneously
- Drug delivery

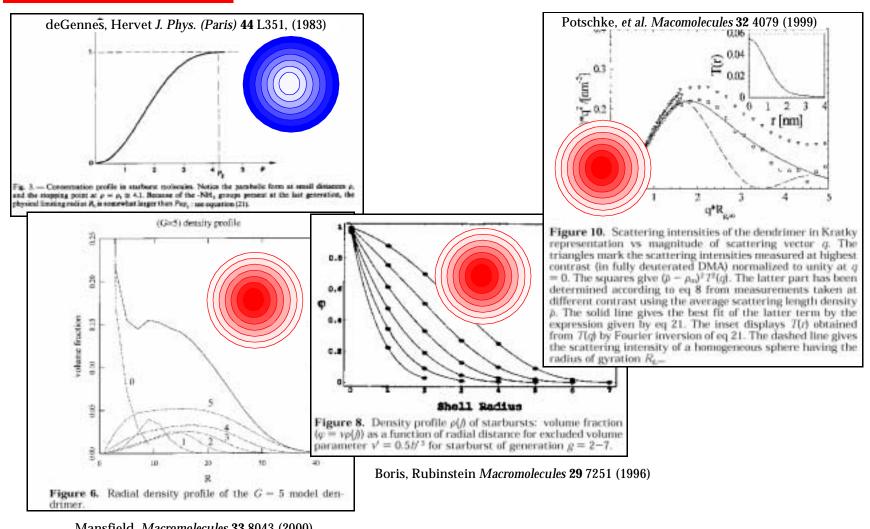
#### Lescanec and Muthukumar

- Short spacer simulation
- Tips dense in center
- Monomers dense in center





#### Other theories, experiments support Filled core



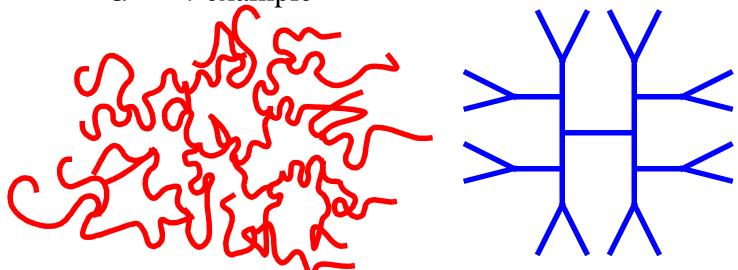
Mansfield, Macromolecules 33 8043 (2000).

# Look at hollow-core model again

a-la Hervet and deGennes:

 $\Box$  G generations, flexible spacers of N monomers

G = 4 example

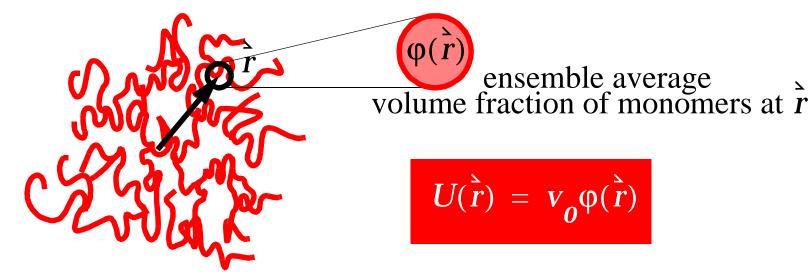


 Excluded volume and chain entropy are the only effects in the Hervet and deGennes calculation

#### • Excluded volume:

2nd virial, mean-field approach:

Energy to insert a monomer at  $\dot{r}$ :  $U(\dot{r})$ 



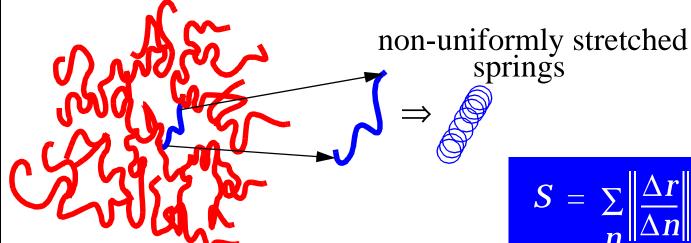
□ Total excluded volume free energy: $E = \sum U(\hat{r}_n)$ 

n

 $\Box$  What is the correct U, or  $\varphi$ ?

#### **Chain entropy**

Gaussian chain segments:



$$S = \sum_{n} \left\| \frac{\Delta r}{\Delta n} \right\|^{2}$$

- Total free energy:  $F = E + S = \sum_{n} \left( \left\| \frac{\Delta r}{\Delta n} \right\|^2 + v_o \phi(r) \right)$ 
  - Self-consistent loop: Find  $\hat{r}(n)$  minimizing F[r], find  $\varphi(r)$ , repeat.

- Further Approximations
- $\Box$  Chemical index *n* and weighting factor f(n):

$$f(n) = 2$$

$$f(n) = 16$$
Smoothed

$$f(n) = 2^{n/G}$$

□ Free energy, saddle point

$$F[r] = \int_{0}^{GN} f(n) \left[ \left\| \frac{dr}{dn} \right\|^{2} + V_{o} \phi(r) \right] dn$$

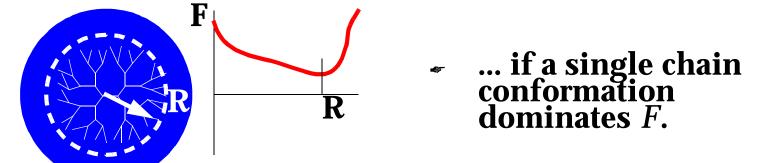
$$-\frac{d}{dn}\left[f(n)\frac{dr}{dn}\right] + f(n)\frac{d\varphi}{dr} = 0 \qquad \Rightarrow \frac{d^2r}{dn^2} - b\frac{dr}{dn} + v_0\frac{d\varphi}{dr} = 0$$

ho Minimizing F gives an ordinary differential eq.

- But, still need  $\varphi(r)$
- Hervet and deGennes make an approximation:

$$\phi \approx \frac{f(n)}{dr/dn}$$
 Multiply number of equivalent chain segments by Monomer density along a single stretched strand BUT, need a unique  $r(n)$ 

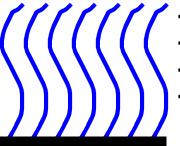
□ **Ok...** 



□ Gives  $\varphi(r)$  growing strongly out to edge of dendrimer.

## Polymer and Dendrimer Brush

# Polymer brush, no branchings:

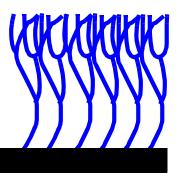


- φis constant
- Tips segregated
- Scaling
- Not selfconsistent

- φ is parabolic
- Ends everywhere
- Self-consistent
- Monodispersity is key constraint
- F is uniform

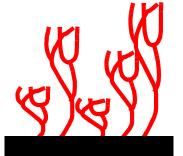


#### Dendrimer brush quite similar:



- φis large at free surface
- Tips segregated
- Scaling
- Not selfconsistent

- φ is still parabolic
- Ends everywhere, concentrated at grafting surface
- Self-consistent
- Monodispersity is key constraint



 $\Box$  Parabolic  $\varphi$ , densest at grafting surface.

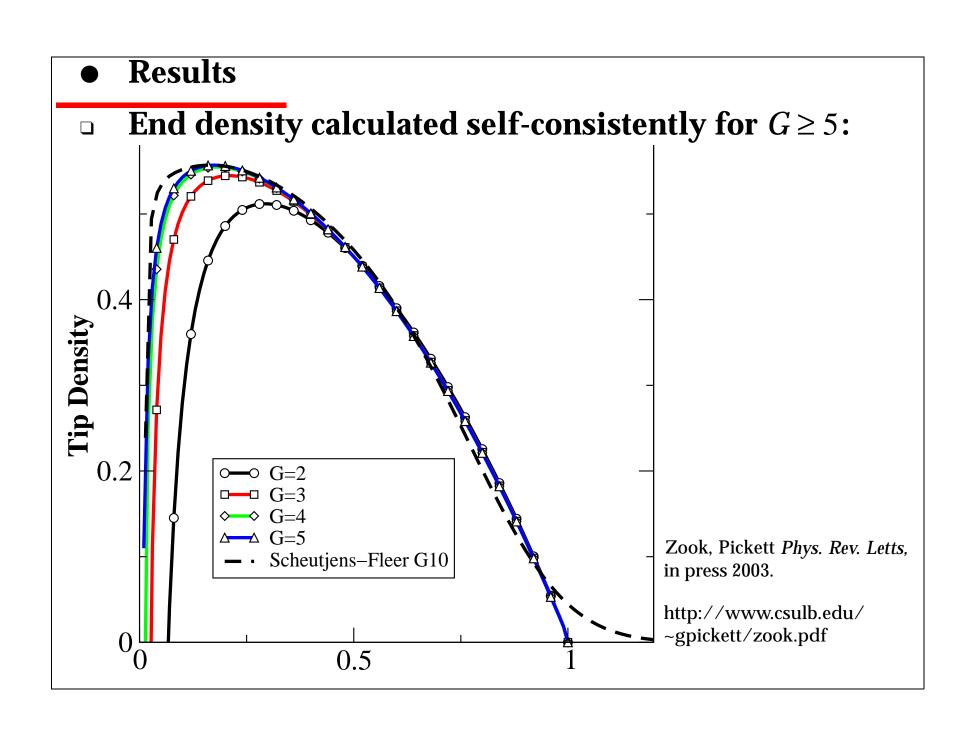
- Parabolic  $\varphi(r)$  is Correct for Dendrimers
- □ 1st order Linear ODE to solve:

$$\frac{d^2r}{dn^2} - b\frac{dr}{dn} + v_o r(n) = 0$$

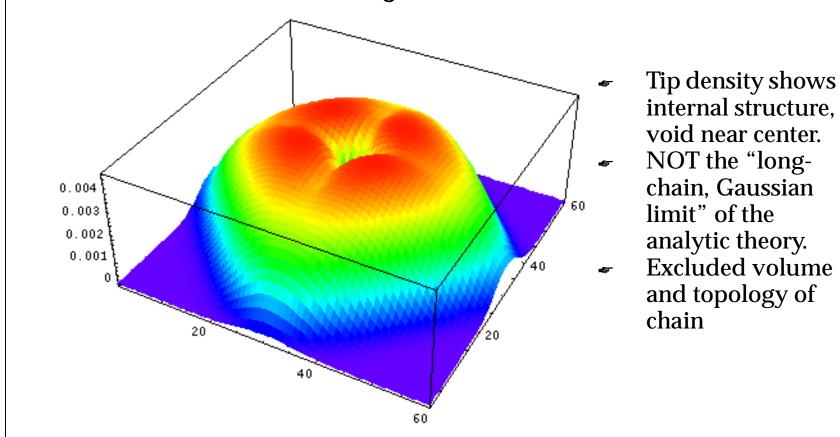
Given r(n) satisfying this ODE, with r(0) = R, r'(0) = 0 can calculate restricted free energy:

$$F[R] = \int_{0}^{GN} f(n) \left[ \left\| \frac{dr}{dn} \right\|^{2} + V_{o} \phi(r) \right] dn = \text{constant}$$

 Dendrimer conformation is a result of many nearly degenerate conformations, spreading the tips from the center out to the edge



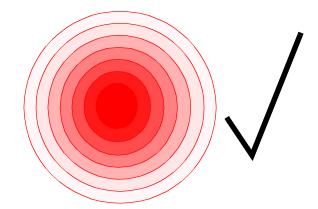
- Interesting Structures Not in Theory
- Short spacers give distinctly non-parabolic density/ density of tips:
- □ N=4, G=8, 2D Scheutjens and Fleer calculation



#### Conclusion

- Hervet and deGennes model predicts Filled Core, not hollow core, when assumptions are relaxed.
- All simulations give filled core.
- Experiments, too.
- Filled core is IT.

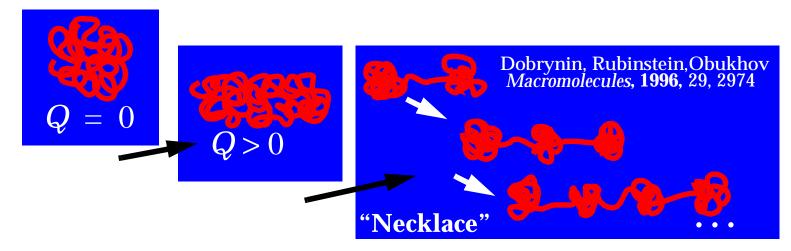




Outline

- Self-consistent field treatment of complex polymer mixtures
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- Conclusions

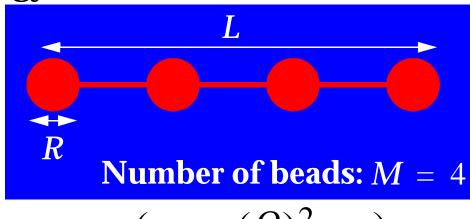
- Polyelectrolyte, poor solvent
- □ Fixed charge  $Q = \alpha N$  on a flexible polymer, N monomers.
- Poor solvent:
  - $\blacksquare$  Q, N control conformation.



Cascade of transitions.

Cartoon theory

Free energy of necklace conformation:



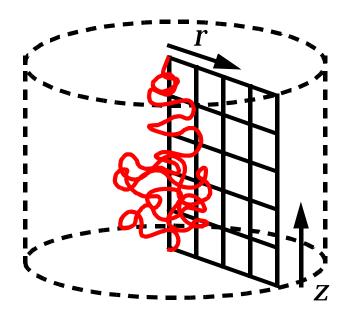
$$F = M\left(\gamma R^2 + \left(\frac{Q}{M}\right)^2 R^{-1}\right) +$$

$$L\gamma + \frac{Q^2}{L}$$

Predicts transitions from 1 to 2 to...

• 2-D Cylindrical lattice

Azimuthal symmetry:

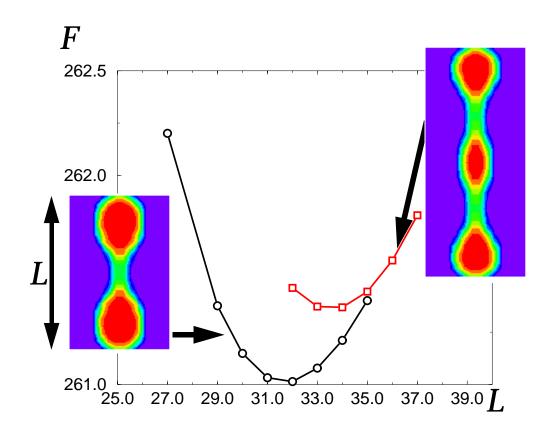


- r,z label annular section of three dimensional space.
- Polymer is held at center of top and bottom plate "bridging"
- Electrostatics, surface energy, chain connectivity are all accounted for

 Variations in 2D, but real 3D structures (highly symmetric).

- Compare structures:
- $f At\ fixed\ N,\ lpha\ vary\ L\ to\ find\ equilibrium\ structures.$

$$N = 250, \alpha = 0.16, \chi = 2.0$$
:

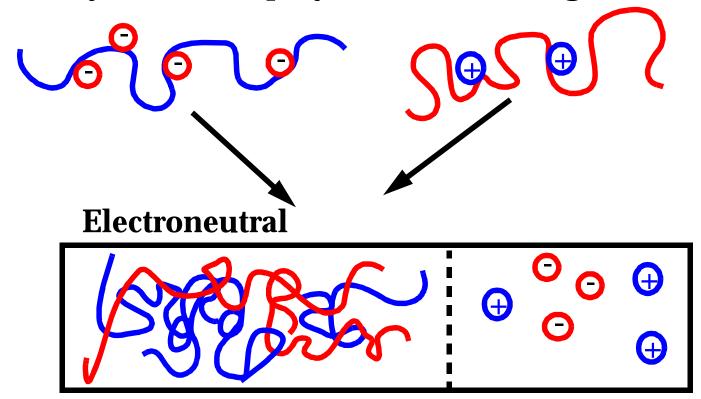


 $L_{eq}$  minimizes F. Possible experiment.

# Diagram of states: **Ω**<sub>8</sub> 0.17 0.16 0.15 0.14 0.13 $\alpha = N^{-1/2}$ N100 400 "Folded" conformations.

• Blend of polyelectrolytes:

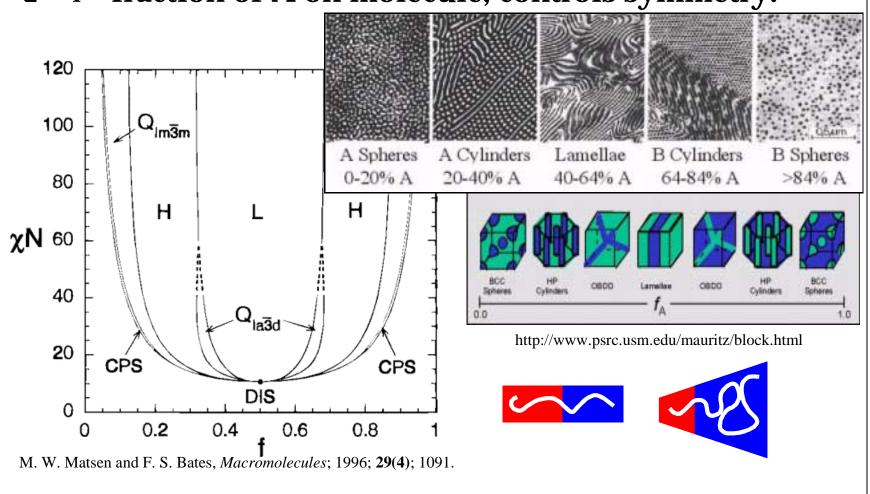
Polycation and polyanions mixed together:



Poly-salt melt... what might it do. Phase separate?

### Asymmetric diblocks

 $\Box$  f= fraction of A on molecule, controls symmetry:



- Coarse-grained Free Energy for Diblocks
- Local interactions

$$F_{\text{local}}[\varphi] = \int \left[\frac{t}{2}\varphi^2 + \frac{k}{2}\nabla\varphi \cdot \nabla\varphi + \varphi^4\right] dx$$

Long-ranged interactions

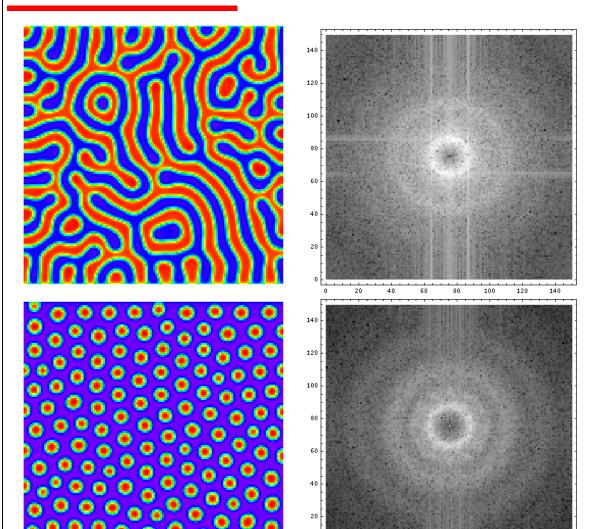
$$F_{\text{long-range}}[\varphi] = \int dx \int dx' B\varphi(x) G(x, x') \varphi(x')$$

Ohta and Kawasaki:

$$\nabla_{\mathbf{X}}^{2}G(\mathbf{X},\mathbf{X}') = -\delta(\mathbf{X}-\mathbf{X}')$$

- Formally, same as electrostatics.
  - A monomers negative, B monomers positive

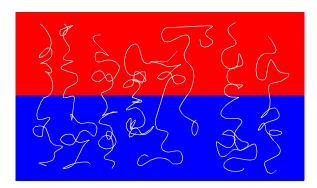
# Minimizing F gives diblock-like structures

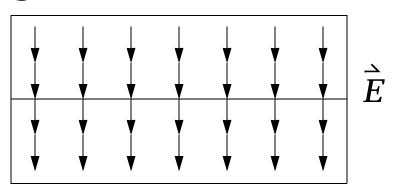


- Cahn-Hilliard dynamics
- Lamellar phase
- Scattering

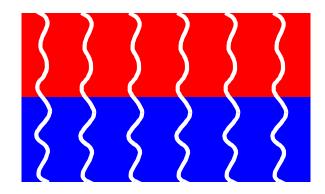
- Cylinder phase
- Scattering

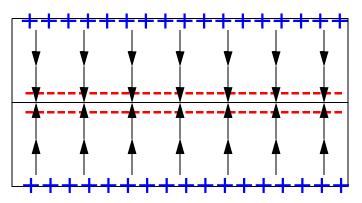
- Electrostatic analogy for Diblocks
- □ Elastic energy ⇔ Electrostatic self-energy
- Semenov, chain stretching similar to electric field:





Alexander, deGennes, and elaborations





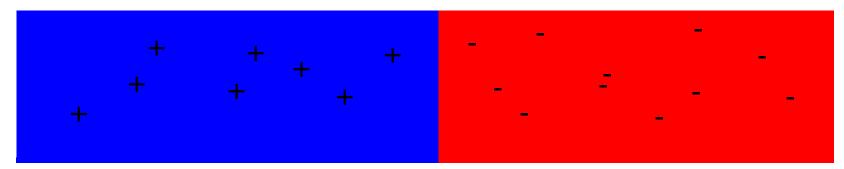
Blend to consider

- Let both chains have the same number of monomers (can be relaxed...)
- Let the CHARGE/monomer on the majority component be fixed.
- Electroneutrality then relates the CHARGE/
   monomer of minority component to composition:

$$0 = \alpha_{\mathbf{A}} f + \alpha_{\mathbf{B}} (1 - f)$$

 Minority chain is more strongly charged than majority chain... synthetic chemistry. Can expect a mesophase.

Phase separation: huge electrostatic costs

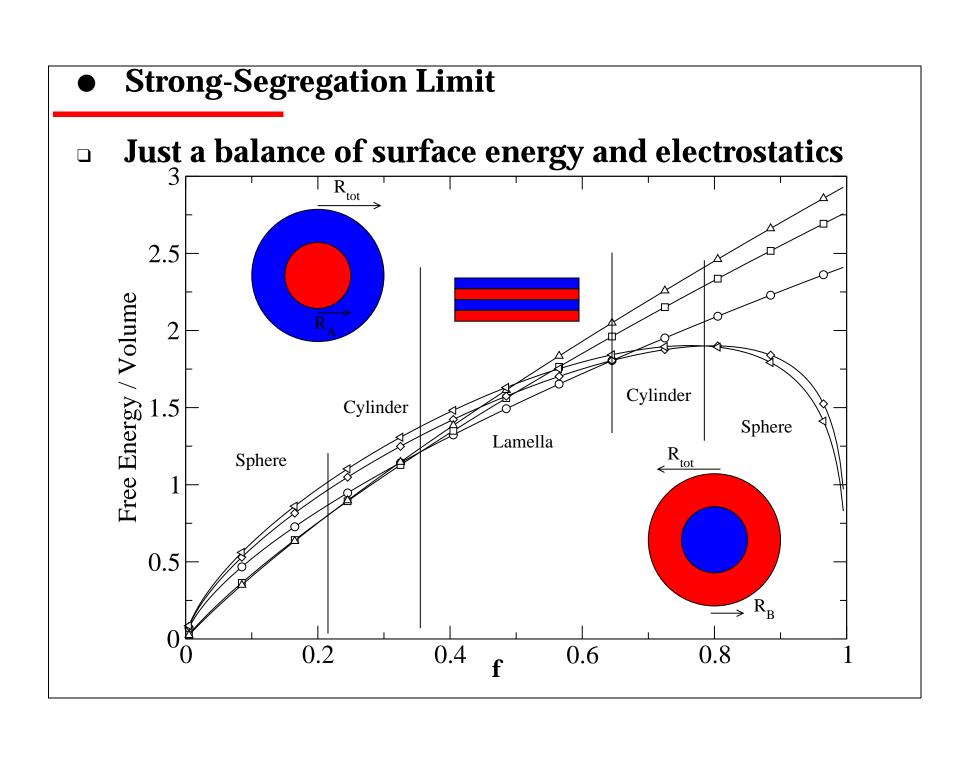


"Collecting like charges"

Single phase: huge specific interactions



N red monomers: total cost  $\chi$  N



• Lattice Electrostatics

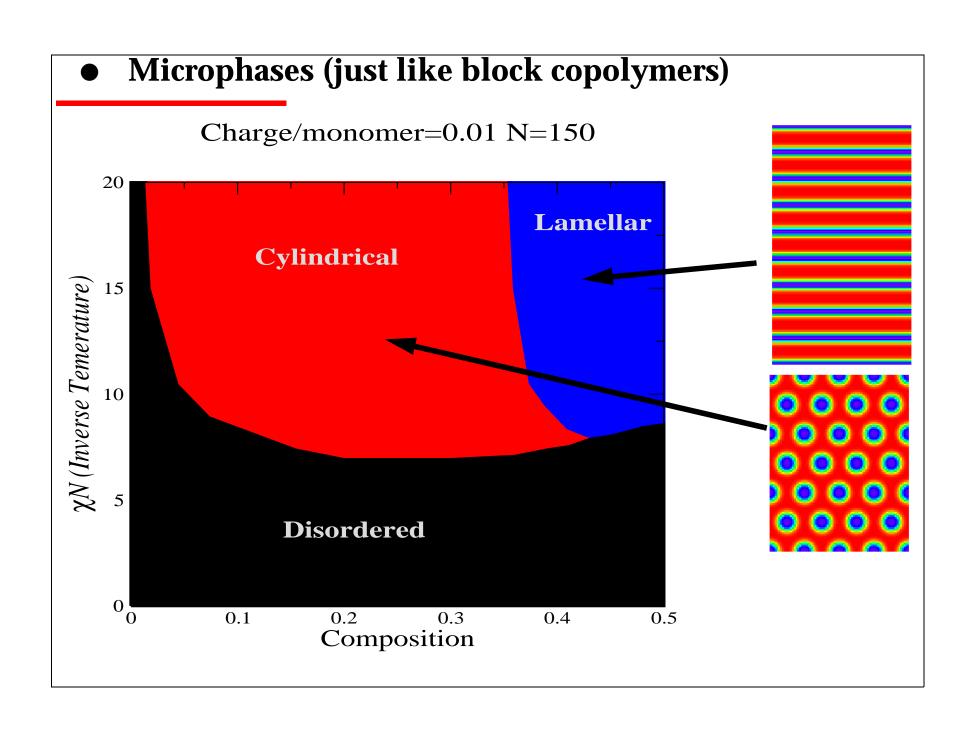
Discretize Laplacian:

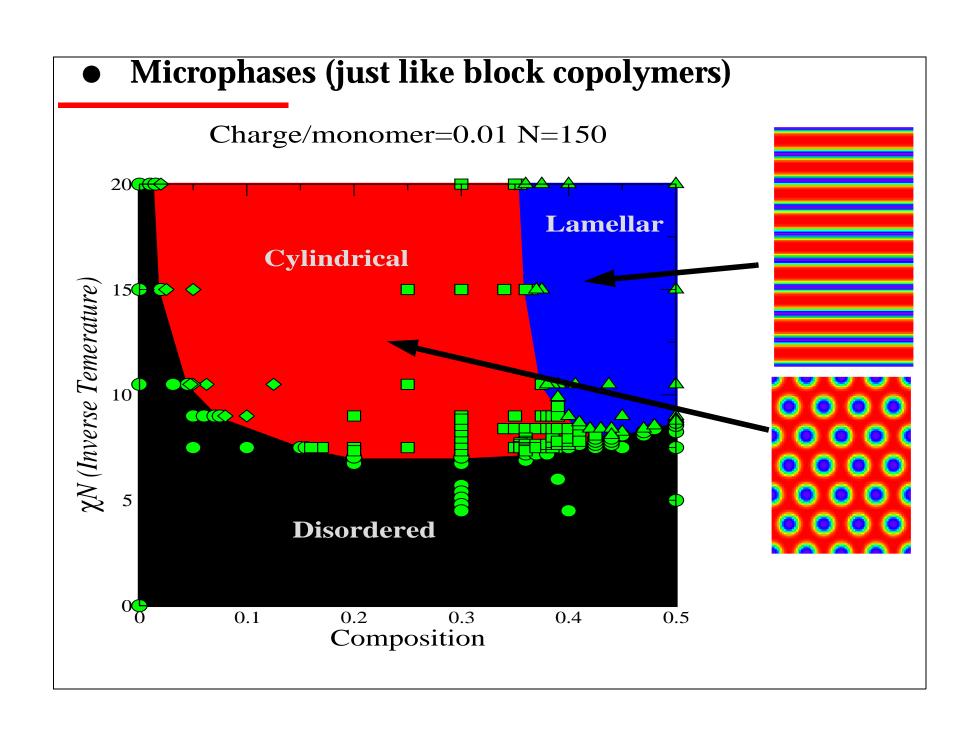
$$\nabla^2 \varphi \Rightarrow \varphi(x, y+1) + \varphi(x, y-1) + \varphi(x+1, y) + \varphi(x-1, y) - 4\varphi(x, y)$$

Gauss' Law discretized:

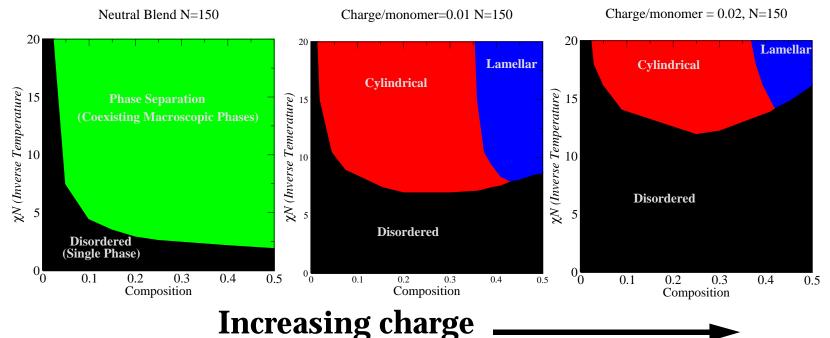
$$\nabla^2 \Phi = 4\pi (\rho_A \phi_A + \rho_B \phi_B)$$

- $\Box$  Solve for  $\Phi$ , electrostatic potential, involves inverting a linear operator on the lattice
- Solved numerically at each iteration by direct inversion.





# Charge compatibilizes the blend



- increasing charge \_\_\_\_
- Simple architectures (just homopolymers) but complex patterns.
- Long-range vs. short-range

# **Films** Lower surface held at a constant potential Upper surface is vacuum Confinement and external field controls morphology **Film** Φ vacuum No Field grounded **External Field** 125 75 100 125

#### Conclusions

#### Dendrimers

 Can't rely on excluded volume and entropy for drug delivery, need something more specific

#### Single charged chains

 Single-chain self-consistent treatment lacks a priori assumptions, points toward clean experiment

# Charged blends

Dynamic control and pattern formation